Stellar winds and coronae from solar-type stars with different metallicities

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Feb. 25th., 2019

Outline

- Solar Wind and Corona by Alfvénic waves
 - Importance & Roles of Density Perturbations
- Active & Inactive Solar Winds
- Stellar wind from low/zero-metallicity solar-type stars
- Solar wind velocity and Magnetic flux tubes

Surface Cartes Surface Surface

HINODE-SOT Solar Wind/Coronal Mass Ejection

SOHO-LASCO

- Energy Source: Fusion reaction at the center
- Surface Convection
- Hot Corona $(T \gtrsim 2 \times 10^6 \text{K})$ & Solar Wind

Fluctuation at the Photosphere

Fluctuation at the Photosphere δv_⊥ in QSs by local correlation tracking



Fluctuation at the Photosphere δv_⊥ in QSs by local correlation tracking





Energetics



Classification of Regions

Withbroe & Noyes (1977)

Region	СН	QS	AR	
LOSS(erg/cm ² s)	8×10^5	3×10^5	107	
Туре	Wind	Cond. & Rad.	Rad.	
CH-Coronal Holes: OS-Quiet Regions:				

CH=Coronal Holes; QS=Quiet Regions;

AR=Active Regions



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The surface convection has sufficient energy.

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- Extract the kinetic energy of the surface convective turbulence
- Lift up the energy to upper layers

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In Situ Heating in the Corona & Wind

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In Situ Heating in the Corona & Wind

Keys:

Energy/Momentum/Mass transfer in the atmosphere T.

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In Situ Heating in the Corona & Wind

Keys:

 Propagation / Reflection / Dissipation of waves

Layers in Atmosphere



Hinode web

- Corona(≳100MK)
- Transition Region (1–100MK)
- Chromosphere (≲10kK)
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Hinode web

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 Huge Density Contrast (16 orders of mag.

from Photosphere to 1au)

Models/Simulations for "Solar Wind" Where in SW ?

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1 Acceleration in inner heliosphere ($\leq 20R_{\odot} \approx 0.1$ au)

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· From Corona Base: Driving

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- 2 Beyond MHD
 - Ioncyclotron resonance ? Axford & McKenzie 1997; Kohl+ 1998
 - Other kinetic effects on the set with set

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- Cool photosphere & chromosphere
 ⇔ Hot corona & wind
- Huge Density Contrast > 15 orders of mag.
- MHD + rad.cooling + conduction



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Forward-type MHD Simulations
- MHD + rad.cooling & thermal conduction

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- Inject Alfvénic fluctuation, $P(v) \propto v^{-1}$, from photosphere

 $(\delta v = 0.7 \text{ km/s with } P = 30 \text{ sec.- } 30 \text{ min.})$

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2D (2.5D)



Matsumoto & Suzuki 2012, ApJ, 749, 8



mesh#: 8,000x 32



Observation & Simulation

Based on Matsumoto & Suzuki (2014)

HINODE/SOT



(1.5D) Suzuki & Inutsuka 2005 ApJ, 632, L49 U





10⁻⁴ 10⁻³ 10⁻² 10⁻¹

10¹

(R-Rs)/Rs

 cool chromosphere / sharp TR / 1MK corona

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- cool chromosphere / sharp TR / 1MK corona
- drive SW with 700-800 km/s





- longitudinal waves in ρ , v_r (slow MHD \approx sound waves) • $\delta B^2 \Rightarrow \delta \rho$
 - $\delta B^2 \Rightarrow \delta \rho$ (Ponderomotive force)

Hollweg 1982; Kudoh & Shibata 1999

• Parametric Decay (Goldstein 1978) in inhomogeneous background

Shoda+ 2018 & Shoda's Talk

⇒ steepening

⇒ shock dissipation

Compressible waves



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transverse waves in v_{\perp}, B_{\perp} (Alfvén waves)

Reflection everywhere

Reflection test

Suzuki & Inutsuka 2006, JGR, A06101

Observation of $\delta \rho$ **by AKATSUKI** AKATSUKI (JAXA): planned as Venus Climate Orbiter



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The 1st orbital injection failed in 2010 Not a satellite but a planet Observe the Sun in June–July, 2011 when AKATSUKI – Sun – Earth along a straight line Detected $\delta \rho$ (\leftrightarrow Sound Waves)



Miyamoto, Imamura et al.2014

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The second orbital injection was successful on Dec.7, 2015

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 $\delta \rho$: a good target for Parker Solar Probe

Suzuki & Inutsuka 2005



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 ~ 90% of the initial Poynting flux reflected back before reaching the corona.

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- ~ 90% of the initial Poynting flux reflected back before reaching the corona.
- The transmitted energy is enough.

Sun in Time – Solar-type stars–



Sun in Time –Solar-type stars– L_X (Güdel 2007)







• Larger $L_X \& \dot{M} \Leftarrow$ fast rotation Skumanich '72; Ayres '97 $L_X \leq 1000 \times L_{X,\odot} \& \dot{M} \leq 100 \times \dot{M}_{\odot}$



Active Young Solar-type Stars:

- Larger L_X & $\dot{M} \leftarrow$ fast rotation skumanich '72; Ayres '97
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Young Suns: Active & Dense Wind but Saturation

$L_{\rm R} - L_{\rm K}$ from Numerical Simulations


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Nonlinear Solar Wind



Nonlinear Solar Wind

If the input energy $\times 2$



The kinetic energy of solar wind ×10!

Nonlinear Solar Wind

If the input energy $\times 2$



The kinetic energy of solar wind ×10! Conservation Law ???

P-gradient ⇒ Change of v_A(= B / √4πρ)⇒Deformation of Wave Shape(=Partial Reflection)



Nonlinear Response of Solar Winds

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Nonlinear Response of Solar Winds $\times 1$ $E_{input} \times 2$ $E_{\text{input}} \times 1$ 1 0.99 0.01 2 0.1 1.9

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Extended Chromosphere in Active Stars Comparing active & present Sun cases

Suzuki, Imada + 2013

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ρ structure



Gas Lifted up by $\delta B^2 \Rightarrow$ Extended Chromosphere

Extended Chromosphere in Active Stars Comparing active & present Sun cases

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Gas Lifted up by $\delta B^2 \Rightarrow$ Extended Chromosphere $\Rightarrow v_{\rm A}$ changes more slowly. \Rightarrow suppression of wave reflection.

Hollweg 1984; Moore et al.1991
• Reflection test

Stellar Winds in a HR diagram

Univ.of California San Diego HP



Massive Main Sequence Stars

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• Line-driven Winds (*P*_{rad} on metallic lines)

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How about Low-mass Main Sequence Stars ?

- Evolution of Elements in the Universe
 - Big Bang: H, He ,Li
 - In Stars: up to Fe + beyond Fe (S-process)
 - Supernovar & Neutron Star Merger: Beyond Fe (R-process)

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First Stars \approx Zero-Metal Stars

- Lifetime of Stars
 - Massive stars have shorter $\tau_{\rm life}$ $\tau_{\rm life} \propto M_{\star}/L_{\star} \sim M_{\star}^{-3}$

Low-mass Zero-metal (Pop.III) Stars $\tau_{\text{life}} \gtrsim \tau_{\text{cosmos}}$ for $M_{\star} \lesssim 0.8 M_{\odot}$.

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Solid objects with ≥ 3km can accrete.Tanikawa+
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- Inject δv by surface convection: $\rho \delta v^3 \propto \sigma T_{\rm eff}^4$
- Equipartition *B* at the photosphere $8\pi p/B^2 = 1 \Rightarrow B \approx kG$
- Super-radially expanding flux tubes
 (B) = Bf = 1.25 G f: filling factor



Tsuneta et al.2008; Shimojo et al.2009; Itoh et al.2010; Shiota et al.2012

Shallower surface convection zone

- Shallower surface convection zone
- Higher luminosity

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Main Sequence Stars with $M_{\star} = 0.7 M_{\odot}$ Yi+ 2001					
$Z(Z_{\odot})$	$R_{\star}(R_{\odot})$	$T_{\rm eff}(K)$	$L(L_{\odot})$	δv_0 (km s $^{-1}$)	<i>B_{r,0}</i> (kG)
1	0.632	4657	0.169	0.641	2.51
0.1	0.620	5576	0.333	0.760	3.05
0.01	0.618	5815	0.391	0.812	3.06
0	0.617	5842	0.397	0.787	3.25

Main Converse Clare with M

Radiative Cooling Function



Cooling Rate (optically thin) $q_{\rm R} = nn_{\rm e}\Lambda \text{ erg cm}^{-3}\text{s}^{-1}$











\dot{M} –Z for Different M_{\star}



 $\dot{M}(Z \le 0.01 Z_{\odot}) \approx (5 - 20) \times \dot{M}(Z_{\odot})$





Poynting flux by Alfvén waves: $L_{\rm A}f = 4\pi r^2 f \rho \langle \delta v^2 \rangle v_{\rm A}$



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Poynting flux by Alfvén waves: $L_{\rm A} f = 4\pi r^2 f \rho \langle \delta v^2 \rangle v_{\rm A}$ Kinetic energy flux: $L_{\rm K,out} = \dot{M} \frac{v_r^2}{2}$ Radiative loss: $(L_{\mathbf{R}}f)_{\mathbf{tr}} =$ $4\pi \int_{r_{t..}}^{r_{out}} dr q_{\rm R} f r^2$

Radiative Loss



Energetics $M_{\star} = 0.7 M_{\odot}$



 $\frac{(L_{\rm A}f)_{\rm tr}}{(L_{\rm A}f)_{\rm 0}} - \frac{L_{\rm K,out}}{(L_{\rm A}f)_{\rm 0}} - \frac{(L_{\rm R}f)_{\rm tr}}{(L_{\rm A}f)_{\rm 0}}$

Summary

- Solar Wind and Corona by Alfvénic waves
 - Density perturbations \Rightarrow Shoda's Talk
- Active & Inactive Solar Winds Small change of energy injection ⇒ Huge difference of SW Reflection triggered global instability _{Suzuki+ (2013)}
- Low/Zero-metallicity solar-type stars drive strong winds
- Solar wind velocity and Magnetic flux tubes: v B/f

Fast/Slow Solar Winds InterPlanetary Scintillation Measurement by STE-lab (Kojima+ 2005)











v_{1au} -f (Wang & Sheeley)



Figure 2 Dependence of velocity on the flux-expansion factor. Triangles are averages of all data in each year. Solid circles are data with weak magnetic-field strength (< 5 G) and small expansion factors (< 100), while open circles are those with strong magnetic-field strength (> 50 G) and a large expansion factor (> 1000). A solid curved line is the best-fit solution for Equation (3), while a dashed line is the average best-fit solution for the 24 years of data from year 1986 to 2009.

 A: Yearly Average

Fujiki et al.2015

● : Weak *B* & Small *f*

○:
 Strong B
 & Large f

v_{1au} - B_{\odot}





Weak B & Small f

0: Strong B & Large f

Figure 4 Dependence of velocity on the photospheric magnetic-field strength. Triangles are averages of all data in each year. Solid circles are data with weak magnetic-field strength (< 5 G) and a small expansion factor (< 100), while open circles are data with a strong magnetic-field strength (> 50 G) and a large expansion factor (> 1000).

 v_{1au} - B_{\odot}/f



Figure 7 Dependence of velocity on the ratio of photospheric magnetic-field strength and a flux-expansion factor. Triangles are averages of all data in a year. Solid circles are SW1 with weak magnetic-field strength (< 5 G) and a small expansion factor (< 100), and open circles are SW2 with strong magnetic-field strength (> 50 G) and a large expansion factor (> 1000). Solid lines are the best-fit regression lines, while dashed lines are the average for the 24 years of data from 1986 to 2009. Fujiki et al.2015

 A: Yearly Average

●: Weak *B* & Small *f*

 O: Strong B & Large f

V–B/f relation

Suzuki 2006, ApJ, 640, L75 Data from Kojima+ (2005)



Energy conservation in a single flux tube gives $v_{1AU} = 300 (\text{km/s})$ $\times \sqrt{5.9 \left(\frac{-\langle \delta B_{\perp} \delta v_{\perp} \rangle_{\odot}}{8.3 \times 10^5 (\text{cm s}^{-1}\text{G})}\right) \left(\frac{B_{r,\odot}(\text{G})}{f_{\text{tot}}}\right) + 3.4 \left(\frac{\gamma}{1.1}\right) \left(\frac{0.1}{\gamma - 1}\right) \left(\frac{T_{\text{C}}}{10^6 (\text{K})}\right) - 4.2}$

V–B/f relation

Suzuki 2006, ApJ, 640, L75 Data from Kojima+ (2005)



Energy conservation in a single flux tube gives

$$v_{1AU} = 300 (km/s) \times \sqrt{5.9 \left(\frac{-\langle \delta B_{\perp} \delta v_{\perp} \rangle_{\odot}}{8.3 \times 10^5 (cm \ s^{-1} \text{G})}\right) \left(\frac{B_{r,\odot}(\text{G})}{f_{\text{tot}}}\right) + 3.4 \left(\frac{\gamma}{1.1}\right) \left(\frac{0.1}{\gamma - 1}\right) \left(\frac{T_{\text{C}}}{10^6 (\text{K})}\right) - 4.2}$$

 $v \leftarrow$ Alfvén waves in super-radially open flux tubes.