70M<sub>の</sub>のブラックホールを持つ

#### とされる連星系LB-1の形成過程

について

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Tanikawa, Kinugawa, Kumamoto, Fujii (2020, PASJ, 72, 39) arXiv:1912.04509

概要

- ・ 天の川銀河の全散開星団内でできるLB-1のような天体の形成率を見積もった。
- ・この形成率からはLB-1の存在を説明できない.
- ・観測,恒星・連星進化理論のどこかに間違いがある と考えられる。

#### Stellar-mass black hole (BH)

- A final state of massive stars
- X-ray binaries and merging BHs
- Not enough information
  - X-ray binaries are short-period binaries, P ≤ 1 day (Corral-Santana et al. 2016).
  - The origin of merging BHs are unknown.





# BHs in long-period binaries

- AS 386: 131 days,  $7M_{\odot}$  compact object (Khokhlov et al. 2018)
- A detached binary in NGC 3201: 167 days,  $4.36M_{\odot}$  compact object (Giesers et al. 2018)
- 2MASS J05215658+4359220: 83 days,  $3.3M_{\odot}$  compact object (Thompson et al. 2019, Science, 366, 637)



$$\frac{M_{\rm CO}^3 \sin^3 i_{\rm orb}}{(M_{\rm giant} + M_{\rm CO})^2} = \frac{K^3 P_{\rm orb}}{2\pi G} (1 - e^2)^{3/2} \sim 0.766 M_{\odot} \rightarrow M_{\rm CO} \gtrsim 2.9 M_{\odot}$$

$$\begin{split} R_{\text{giant}} &= v_{\text{spin}} P_{\text{spin}} / 2\pi \sim \frac{23 \pm 1 R_{\odot}}{\sin i_{\text{spin}}} \left( \frac{v_{\text{spin}}}{14.1 \text{kms}^{-1}} \right) \left( \frac{P_{\text{spin}}}{82.2 \text{day}} \right) \\ M_{\text{giant}} &= g_{\text{giant}} R_{\text{giant}}^2 / G \sim \frac{4.4_{-1.5}^{+2.2} M_{\odot}}{\sin^2 i_{\text{spin}}} \left( \frac{R_{\text{giant}}}{23 R_{\odot}} \right)^2 \left( \frac{g}{10^{2.35} \text{cms}^{-2}} \right) \\ P_{\text{spin}} \sim P_{\text{orb}}, e \sim 0 \rightarrow i_{\text{spin}} \sim i_{\text{orb}} \sim i \text{ (synchronized)} \end{split}$$

#### LB-1

- $8M_{\odot}$  B-type star  $70M_{\odot}$  BH
- $a \sim 1$  au,  $e \sim 0.03$ ,  $Z \sim Z_{\odot}$
- L, T, and g constrain B-type star mass.
- The ratio of radial velocity determines BH mass.





### What's surprising?

- High metallicity  $(Z \sim Z_{\odot})$ 
  - Stellar wind mass loss reduces BH mass to  $\leq 20M_{\odot}$ .
  - The mass loss rate should be 5 times smaller than previously thought.
- Circular orbit ( $e \sim 0.03$ )
  - Circularization timescale  $(\sim 10^{14} \text{ yr})$  is much more than the Hubble time.



#### Reduced stellar wind

- BH progenitors should have  $M_{\rm tot} \gtrsim 70 M_{\odot}$ and  $M_{\rm c,He} \lesssim 45 M_{\odot}$ .
  - BH progenitors with  $M_{\rm c,He} \gtrsim 45 M_{\odot}$ reduce BH masses to  $M_{\rm BH} \sim 45 M_{\odot}$ through mass loss of pulsatoinal pair instability (PPI).
  - BH progenitors with  $M_{c,He} \gtrsim 65 M_{\odot}$ leave no remnants due to pair instability supernovae (PISNe)
  - GW observation supports PPI/PISN.
- The binary size  $(a \sim 1au \sim 200R_{\odot})$ 
  - $a_i \lesssim 1$ au ... MS merger
  - $a_i \gtrsim 1$ au ... Common envelope
  - $a_i \gg 1$ au ... No interaction ( $a \gg 1$ au)



# Possibility of double BHs

- The merger time through gravitational wave is  $\sim 10^4$  yr.
- The merger time is smaller than the lifetime of the B-type star by three order of magnitude.
- This probability is quite low.
- (Shen et al. 2019)



#### Possible scenario

- Binary evolution
- Hierarchical multiplicity
  - Inner BH-BHs
  - (More complicating channels)
- Dense stellar cluster
  - Capture of a B-type star by a BH
  - More complicating channels

# Counter opinions on " $70M_{\odot}$ "

- No evidence that  $H_{\alpha}$  is associated with the BH
- Radial velocity variability disappears when  $H_{\alpha}$  absorption by the B-type star is considered.
- $H_{\alpha}$  may be associated with circumbinary materials.



El-Badry, Quataert (2020; see also Abdul-Masih et al. 2019)



phase

# "Postgenitor" problem

- LB-1 system will evolve to an ULX source in future.
  - Roche-lobe overflow will starts when the B-type star enters into a Hertzsprung gap (HG) phase.
  - The HG star rapidly expands, and achieves a high accretion rate onto the BH.
- The number of ULXs inferred by LB-1 is larger than observed in the MW by an order of magnitude.



#### Our stance

- The presence of the  $70M_{\odot}$  BH may be doubtful.
- However, another theoretically-challenging binary may be reported in future.
- The usual meaning of the "theoretically-challenging" is "challenging in the framework in binary evolution".
- We use this opportunity to notice dynamical formation of a binary in dense stellar clusters, using LB-1 as a good example.

# The most efficient process

- 1. Collision of a naked He star with a MS star which has a B-type companion.
  - The He star must not have Hydrogen envelope.
- 2. The collision product and B-type companion form a binary system.
- 3. The binary system is circularized through dynamical tide of the collision product's envelope.
- 4. The collision product collapses to a  $70M_{\odot}$  BH.
  - It can avoid PPI/PISN because of small He core.

In an open cluster of the MW galaxy



# Collision rate

• Formation rate of PI-gap BHs in OCs

• 
$$\dot{N}_{\rm PIgap} \sim 2 \times 10^{-6} \left( \frac{f_{\rm PIgap}}{0.002} \right) \left( \frac{\rho_{\rm oc}}{10^4 M_{\odot} {\rm pc}^{-3}} \right) \left( \frac{\eta_{20}}{0.003 M_{\odot}^{-1}} \right) \left( \frac{f_{\rm oc}}{0.2} \right) \left( \frac{\dot{M}_{\rm mw}}{2M_{\odot} {\rm yr}^{-1}} \right) [{\rm yr}^{-1}]$$

• Formation path fraction

• 
$$\frac{\Gamma_{\rm nHe}}{\Gamma_{\rm eHe}} \sim 10^{-2} \left(\frac{N_{1,\rm nHe}/N_{1,\rm eHe}}{2}\right) \left(\frac{M_{12,\rm nHe}/M_{12,\rm eHe}}{0.7}\right) \left(\frac{R_{12,\rm nHe}/R_{12,\rm eHe}}{0.01}\right)$$

• Collision rate

• 
$$\dot{N}_{\text{coll}} = \dot{N}_{\text{PIgap}} \frac{\Gamma_{\text{nHe}}}{\Gamma_{\text{eHe}}} P_{\text{b}} \sim 3 \times 10^{-9} \left(\frac{\dot{N}_{\text{PIgap}}}{2 \times 10^{-6} \text{ yr}^{-1}}\right) \left(\frac{\Gamma_{\text{nHe}}/\Gamma_{\text{eHe}}}{10^{-2}}\right) \left(\frac{P_{\text{b}}}{0.1}\right) \text{ [yr}^{-1}$$



### Circularization

- The binary is rapidly circularized through tidal interaction.
- If the collision product collapses to a BH before swallowing the B-type star, the binary becomes LB-1.
  - The collapse time is at random, since the naked He star wandered in an OC for a long time.
- Circularization time

.

$$t_{\rm cric} \sim 2 \times 10^4 \left(\frac{R_{\rm coll}}{100R_{\odot}}\right)^{-9} \text{[yr]}$$

• Kelvin-Helmholtz time (expansion time)

• 
$$t_{\rm KH} \sim 2 \times 10^4 \left(\frac{M_{\rm coll}}{70M_{\odot}}\right)^2 \left(\frac{R_{\rm coll}}{100R_{\odot}}\right)^{-1} \left(\frac{L_{\rm coll}}{10^5 L_{\odot}}\right)^{-1} \, [\rm yr]$$

• Surviving probability

$$P_{\text{surv}} = t_{\text{KH}} / t_{\text{coll,life,max}} \sim 0.1 \left(\frac{t_{\text{coll,life,max}}}{0.2 \text{Myr}}\right)^{-1}$$



# The formation rate



• The number of LB-1-like systems in the MW

$$N_{\rm LB1} \sim 0.01 \left( \frac{\dot{N}_{\rm coll}}{3 \times 10^{-9} {\rm yr}^{-1}} \right) \left( \frac{P_{\rm surv}}{0.1} \right) \left( \frac{T_{\rm B}}{40 {\rm Myr}} \right)$$

• No chance to form LB-1-like systems in OCs

### Other stellar collisions

- Collision of He stars with H envelope does not work.
  - He star have  $R \gg a$ .
- Collision products of two MSs or two naked He stars cannot avoid PPI/PISN
- Collision rate of BH and other stars is lower than or similar to the above process.

	MS	He star	Naked He star	BH
MS	PPI/ PISN			
He star	R>>a	R>>a		
Naked He star	Done	R>>a	PPI/ PISN	
BH	Similar rate	R>>a	Lower rate	Lower rate

#### Possible scenario

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- Hierarchical multiplicity
  - Inner BH-BHs
  - (More complicating channels)
- Dense stellar cluster
  - Capture of a B-type star by a BH
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# Summary

- A  $70M_{\odot}$  BH in LB-1 has been reported.
- The presence may be doubtful, but is under dispute.
- We have examined the formation rate of LB-1, but LB-1 has no chance to be formed through dynamical interactions, and hierarchical triple systems, if the standard model of single and binary stars is correct.
- We don't deny the presence of  $70M_{\odot}$  BHs in wide binaries with  $\gg 200R_{\odot}$  under metal-poor environments.

# Back-up slides

### What is LB-1 in reality?

- The B-type star can be a stripped helium star with  $\sim 1.1 M_{\odot}$  (Irrgang et al. 2019).
- The luminosity is consistent if the Gaia distance is adopted (Eldridge et al. 2019; Irrgang et al. 2019).
- The unseen companion can be a neutron star.



# Capture processes

- At first, there is no circularization process
- OC-origin:  $N_{\rm b} \sim 0.7 \rightarrow N_{\rm b,cir} \sim 7 \times 10^{-4}$
- GC-origin: No B-type star
- Interstellar space-origin:  $N_{\rm b} \sim 7 \times 10^{-8}$



# Hierarchical triple (1)

- The merger product should be  $\gtrsim 70 M_{\odot}$ .
- If it has a radius of  $\gg 200M_{\odot}$ , it is a He star.
- It experiences common envelope evolution with the B-type star.
  - It loses its envelope, and collapses to  $a \leq 45 M_{\odot}$  BH.
  - It merges with the B-type star, and the system should not be a binary system.
- The inner binary should be separated from the B-type star by  $\sim 200R_{\odot}$ , and never has no interaction with the B-type star.



# Hierarchical triple (2)

- The separation of the inner binary should be  $\leq 100R_{\odot}$ . Otherwise, the system is unstable (Harrington 1972; Mardling, Aarseth 1999).
- The primary star of the inner binary should be  $\gtrsim 35M_{\odot}$ .
  - $\leq 100 M_{\odot}$  stars exceed ~  $100 R_{\odot}$ when they are in Hertzsprung gap phases. The inner binary experiences a Case B merger.
  - $\gtrsim 100 M_{\odot}$  stars exceed ~  $100 R_{\odot}$ when they are in MS phases. The inner binary experiences a MS-MS merger. The merger product cannot avoid PPI/PISN.



#### Case B merger

- When the primary star is in a Hertzsprung gap phase, the binary can experience Case B merger.
- But, the merger product has  $\sim 200R_{\odot}$ , and merges with the B-type companion.
  - A  $35M_{\odot} + 35M_{\odot}$  merger product gets the smallest radius  $\gtrsim 200R_{\odot}$ .
  - The mass ratio of the merger product to the B-type star is high  $\gtrsim 10$ .



log(T) (K)

Justham et al. (2014)

#### Pulsational Pair Instabiliity



**Fig. 3.** Adopted models for pair-instability pulsation supernova mass loss. For a given He core mass we show the mass of a star after PPSN mass loss. Moderate PPSN mass loss is adopted directly from Leung et al. (2019), while its modified (50% reduced mass loss) version is presented as weak PPSN model. Strong PPSN are adopted from Belczynski et al. (2016c).