

ブラックホール連星形成の理論的研究
特に全質量 $150M_{\odot}$ であるGW190521を
初代星連星で作れるかどうかについて

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- Tanikawa et al. (2020a, MNRAS, 495, 4170)
- Tanikawa et al. (2020b, arXiv:2008.01890)
- Tanikawa et al. (2020c, arXiv:2010.07616)

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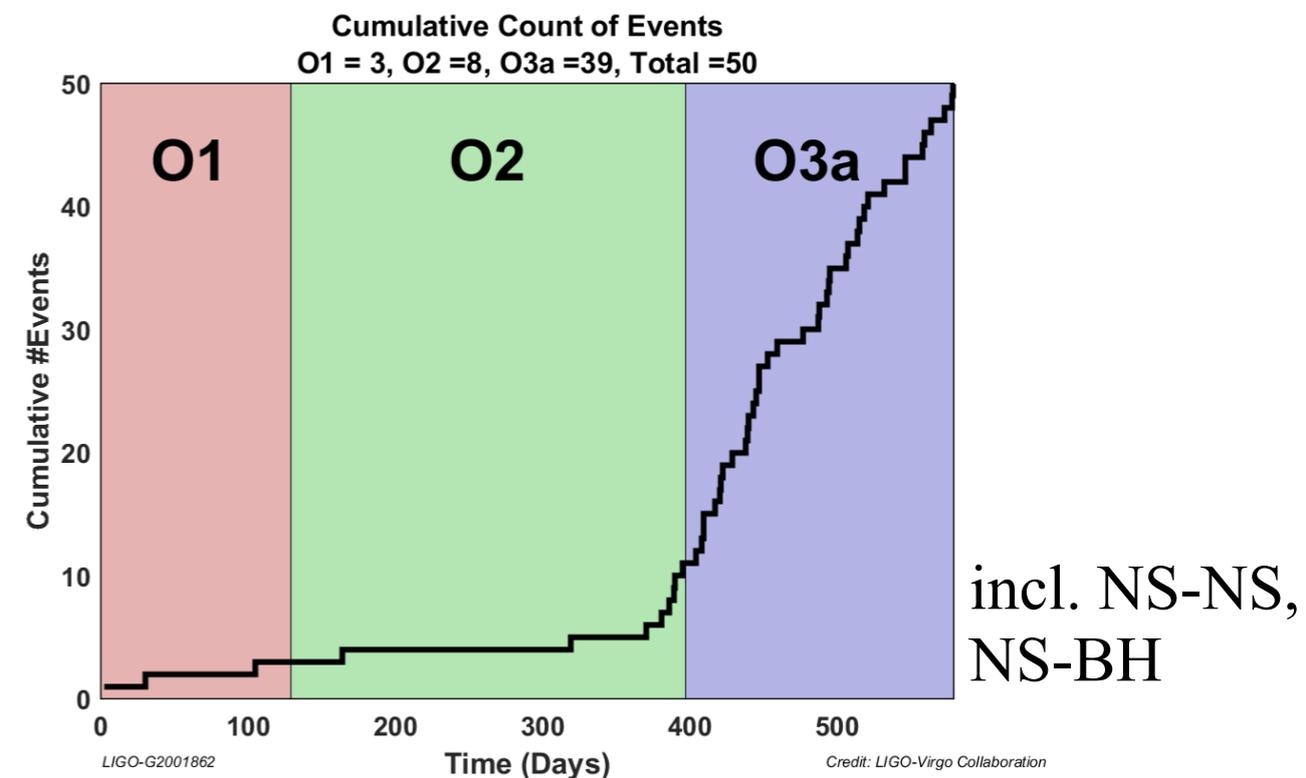
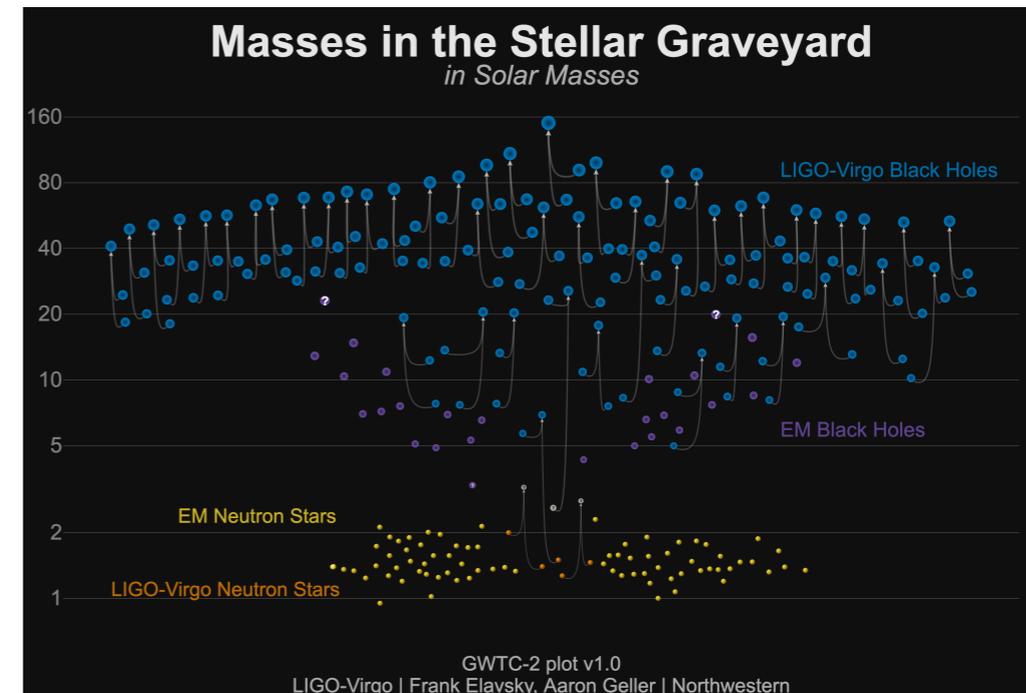
- Pair instability (PI) mass gap and GW190521
- Difficulty to form PI mass gap BH
- Pop. III binary scenarios
- Our study

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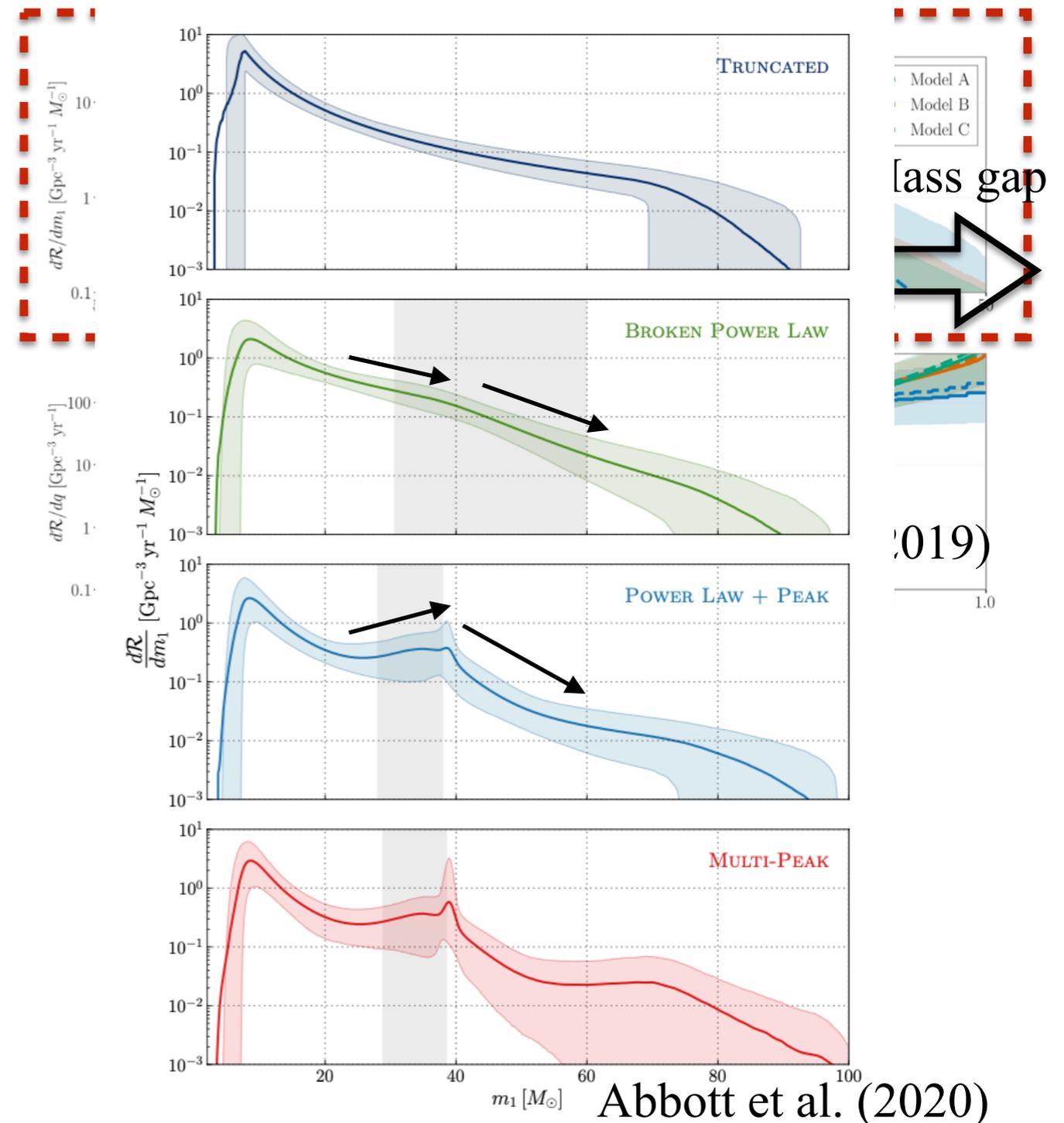
Merger of binary black holes

- Gravitational wave (GW) observations have rapidly increased the number of discovered BH-BHs.
- The first detection is 2015 (GW150914).
- 10 BH-BHs were found until 2017 (O1/O2)
- 44 BH-BHs have been found until now (O1/O2/O3a).
- The number of BHs is larger than that discovered by X-ray observations.



BH mass distribution

- The BH mass distribution appears not to have BHs with $\gtrsim 50M_{\odot}$.
- Second mass gap
- Higher mass gap
- Pair instability (PI) mass gap
- No BH with $50 - 130M_{\odot}$ due to pulsational PI (PPI) and PI supernova (PISN)?



PPI and PISN

- Pulsational Pair Instability (PPI)

- $40 \lesssim M_{c,He,preSN}/M_{\odot} \lesssim 60$

- He core partially disrupted

- $M_{bh} \sim 40M_{\odot}$

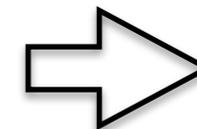
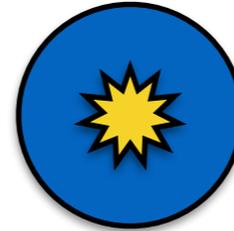
- Pair instability supernova (PISN)

- $60 \lesssim M_{c,He,preSN}/M_{\odot} \lesssim 130$

- He core completely disrupted

- No remnant

$$40 \lesssim M_{c,He,preSN}/M_{\odot} \lesssim 60$$

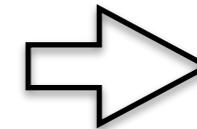
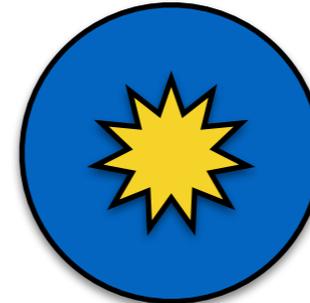


$$M_{bh}/M_{\odot} \sim 40$$



Partial disruption

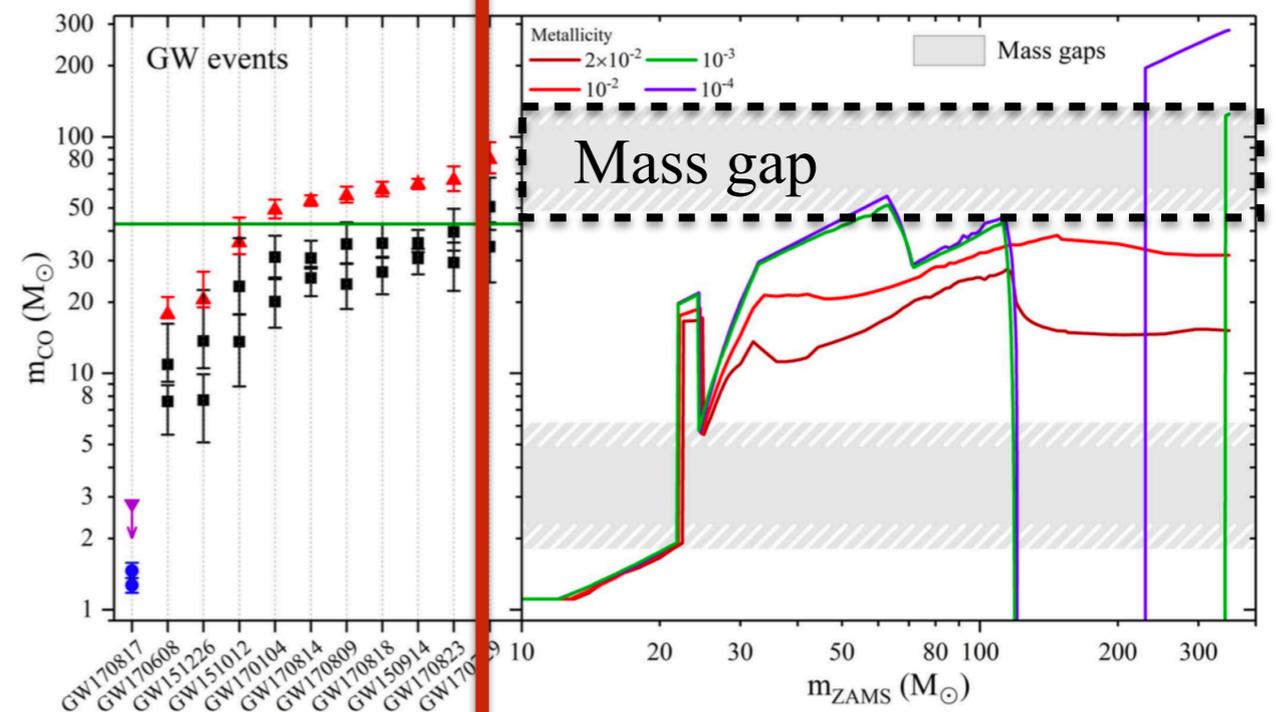
$$60 \lesssim M_{c,He,preSN}/M_{\odot} \lesssim 130$$



No remnant



Complete disruption



Abbott et al. (2019)

GW190521

- Merger of $85_{-14}^{+21} M_{\odot}$ and $66_{-18}^{+17} M_{\odot}$ BHs
- The primary BH has only a 0.32% probability of being below $65 M_{\odot}$.
- At least one BH lies within the PI mass gap.

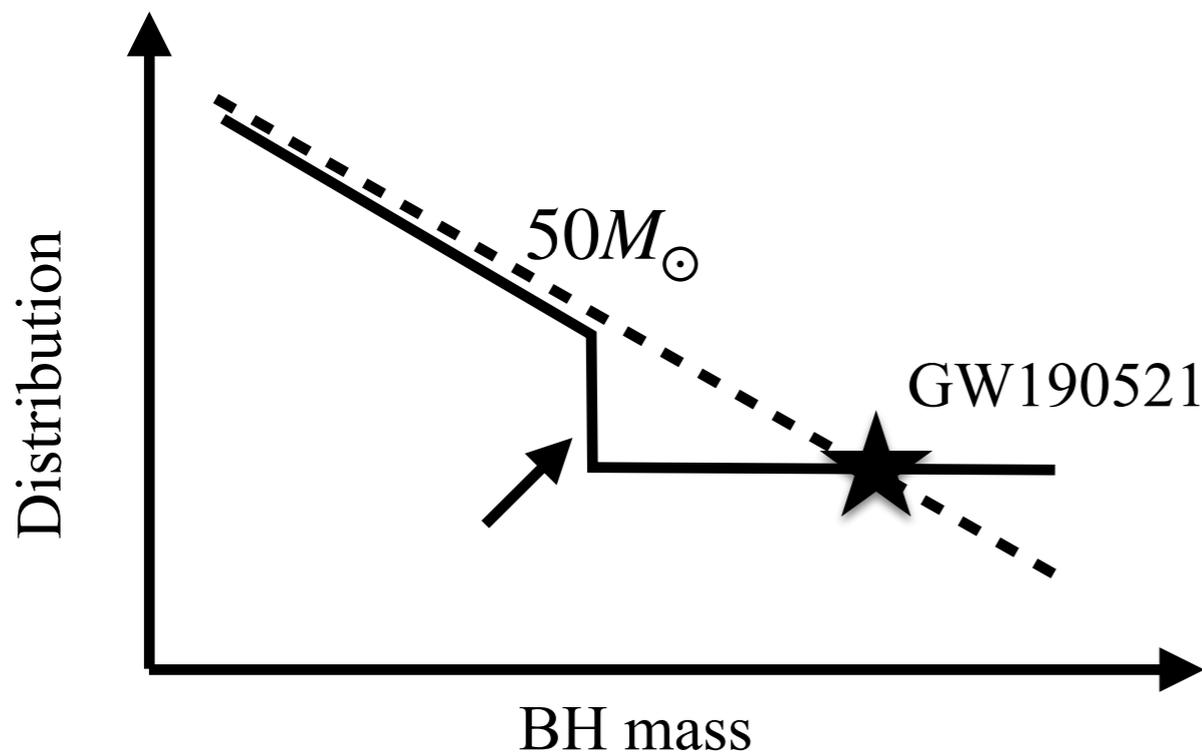
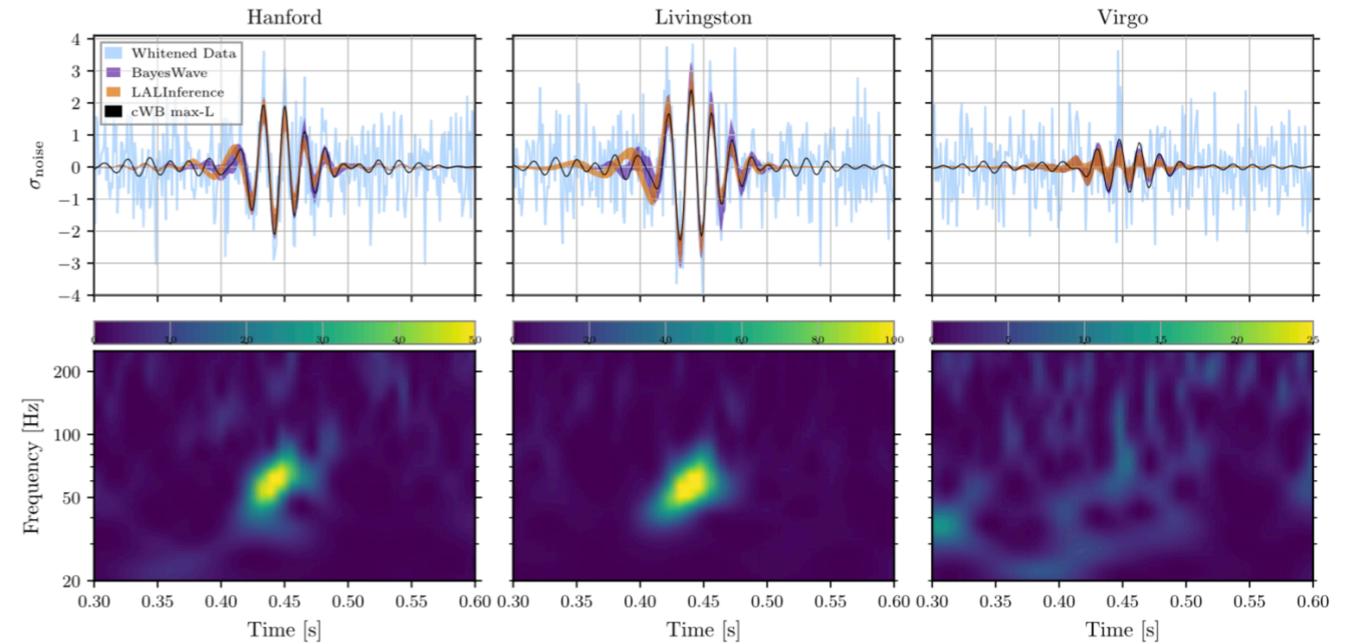


TABLE I. Parameters of GW190521 according to the NRSur7dq4 waveform model. We quote median values with 90% credible intervals that include statistical errors.

Parameter	
Primary mass	$85_{-14}^{+21} M_{\odot}$
Secondary mass	$66_{-18}^{+17} M_{\odot}$
Primary spin magnitude	$0.69_{-0.62}^{+0.27}$
Secondary spin magnitude	$0.73_{-0.64}^{+0.24}$
Total mass	$150_{-17}^{+29} M_{\odot}$
Mass ratio ($m_2/m_1 \leq 1$)	$0.79_{-0.29}^{+0.19}$
Effective inspiral spin parameter (χ_{eff})	$0.08_{-0.36}^{+0.27}$
Effective precession spin parameter (χ_p)	$0.68_{-0.37}^{+0.25}$
Luminosity Distance	$5.3_{-2.6}^{+2.4} \text{ Gpc}$
Redshift	$0.82_{-0.34}^{+0.28}$
Final mass	$142_{-16}^{+28} M_{\odot}$
Final spin	$0.72_{-0.12}^{+0.09}$
$P(m_1 < 65 M_{\odot})$	0.32%
\log_{10} Bayes factor for orbital precession	$1.06_{-0.06}^{+0.06}$
\log_{10} Bayes factor for nonzero spins	$0.92_{-0.06}^{+0.06}$
\log_{10} Bayes factor for higher harmonics	$-0.38_{-0.06}^{+0.06}$

Abbott et al. (2020)

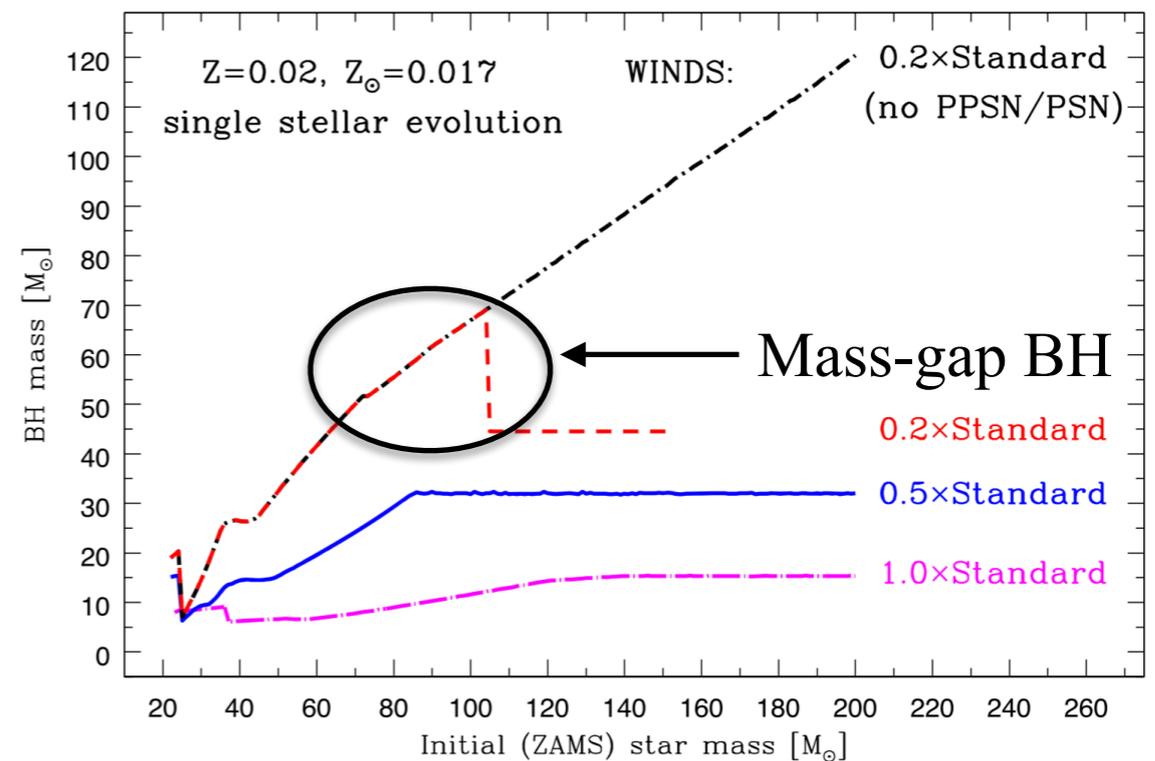
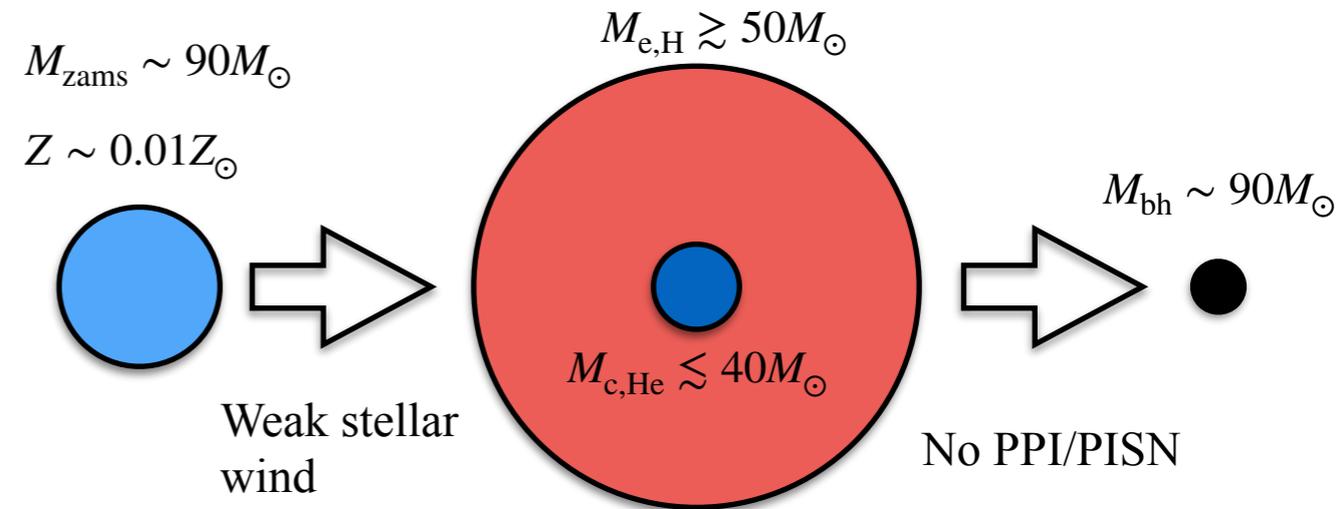
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Single star evolution

- It is not hard to form mass-gap BH through single star evolution.
- Formation Process
 - A star with $M_{\text{zams}} \sim 90M_{\odot}$ and $Z \sim 0.01Z_{\odot}$.
 - Evolution to a BH progenitor with $M_{\text{tot}} \sim 90M_{\odot}$ and $M_{\text{c,He}} \lesssim 40M_{\odot}$.
 - Collapse to $\sim 90M_{\odot}$ BH without PPI/PISN owing to small He core mass.

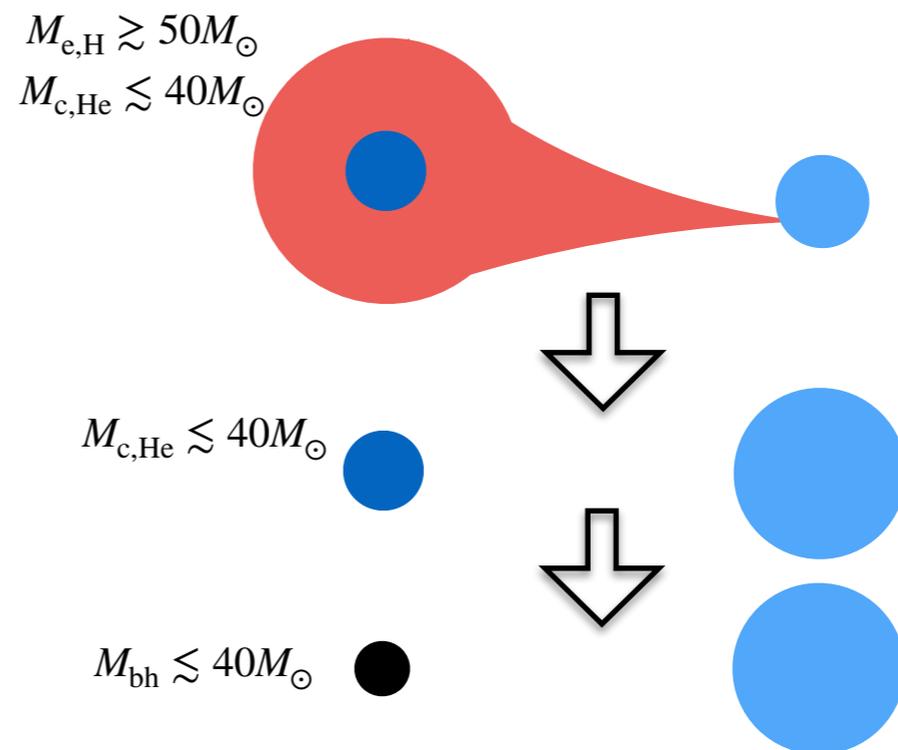
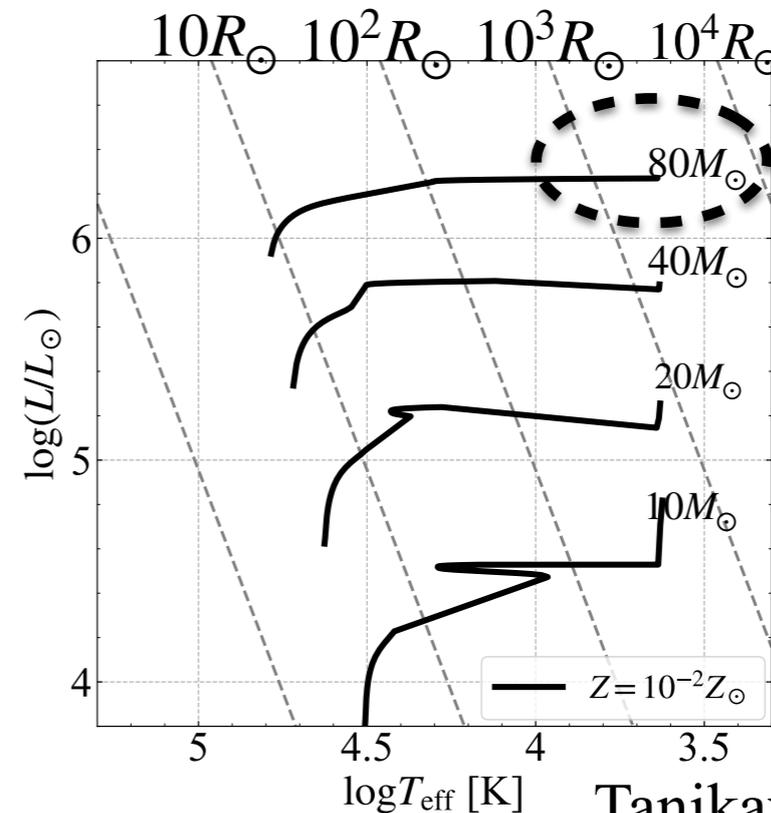
• Light He core, massive H envelope



Belczynski et al. (2020)

Binary star evolution

- Merger of $85M_{\odot}$ and $66M_{\odot}$ BHs
- Merger time $\lesssim 10$ Gyr
- $a \lesssim 10^2 R_{\odot}$, $e \sim 0$
- A star with $M_{\text{zams}} \gtrsim 80M_{\odot}$ expands to $R \gtrsim 10^3 R_{\odot}$.
- The star loses its H envelope, stripped by its companion star.
- No massive H envelope, no mass-gap BH.



Many scenarios other than binary evolution

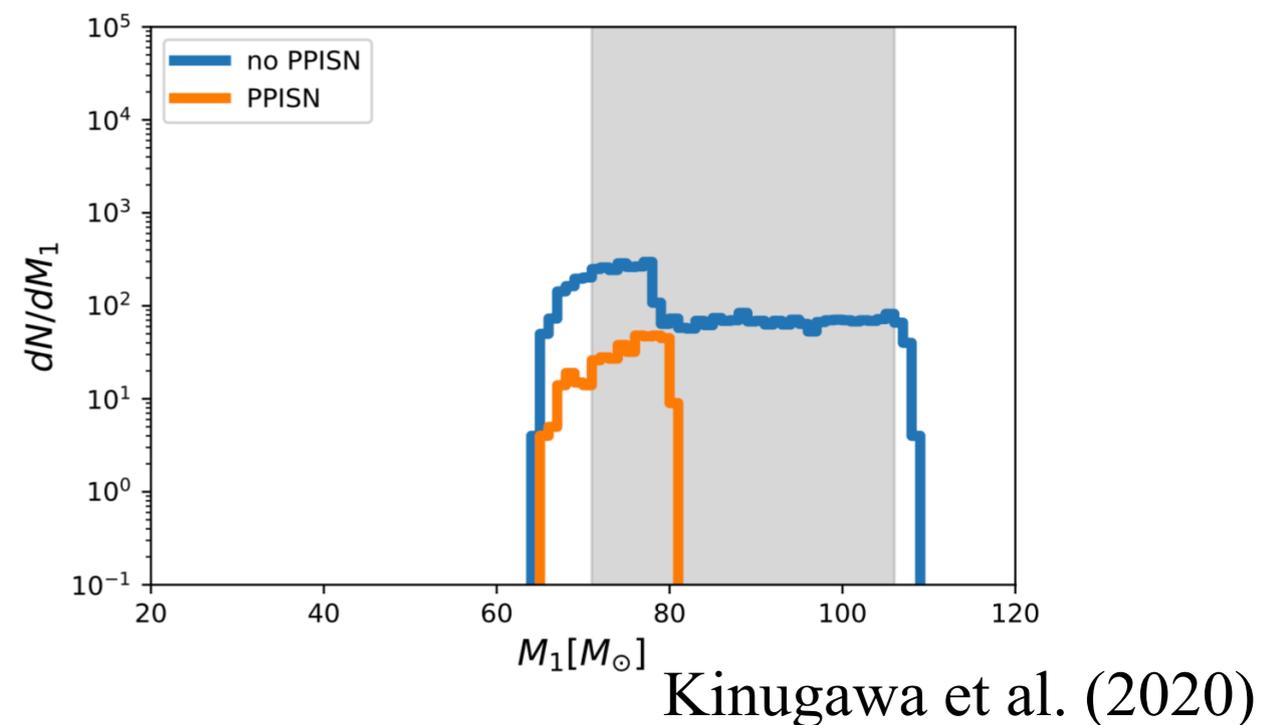
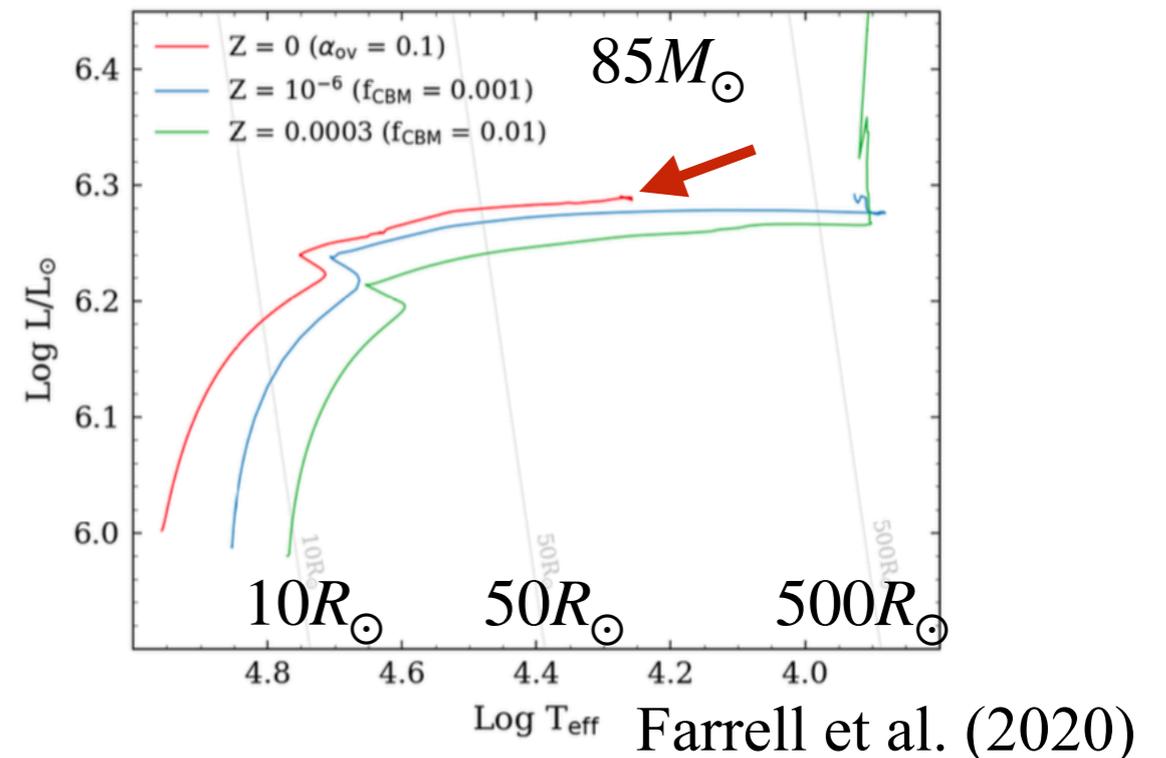
- Globular clusters, open clusters, AGN disks (e.g. Rodriguez et al. 2019; Di Carlo et al. 2020; Yang et al. 2019)
- PPI/PISN occurs in $M_{\text{c,He}} \gtrsim 90M_{\odot}$ if the $^{12}\text{C}(\alpha, \gamma)^{16}\text{O}$ rate is 3 times smaller than the standard one (Takahashi 2018; Farmer et al. 2020; Costa et al. 2020).
- Many more ...
- Our study of GW190521 formation
 - Binary evolution isolated from star clusters
 - PPI/PISN for $M_{\text{c,He}} \gtrsim 40M_{\odot}$ as usual

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Pop. III binary star evolution

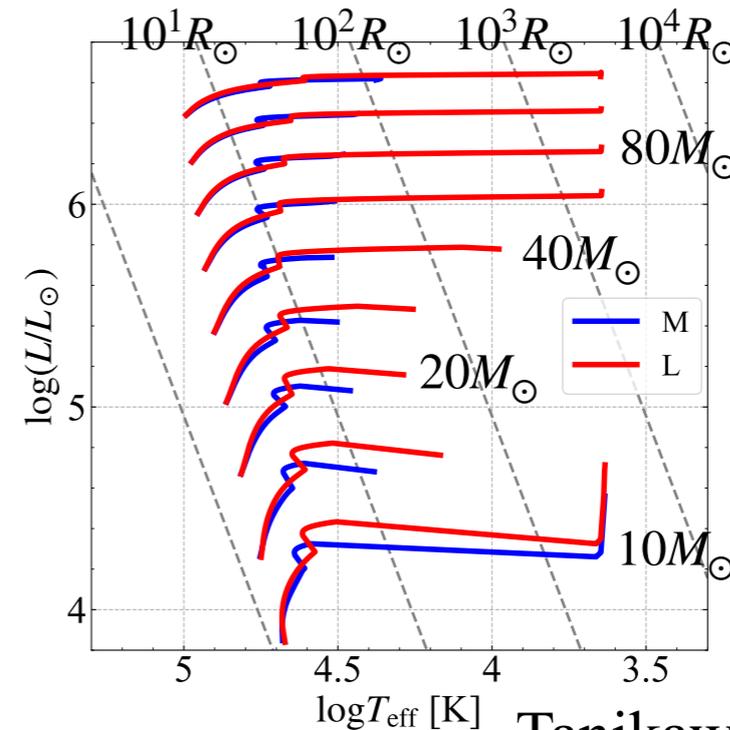
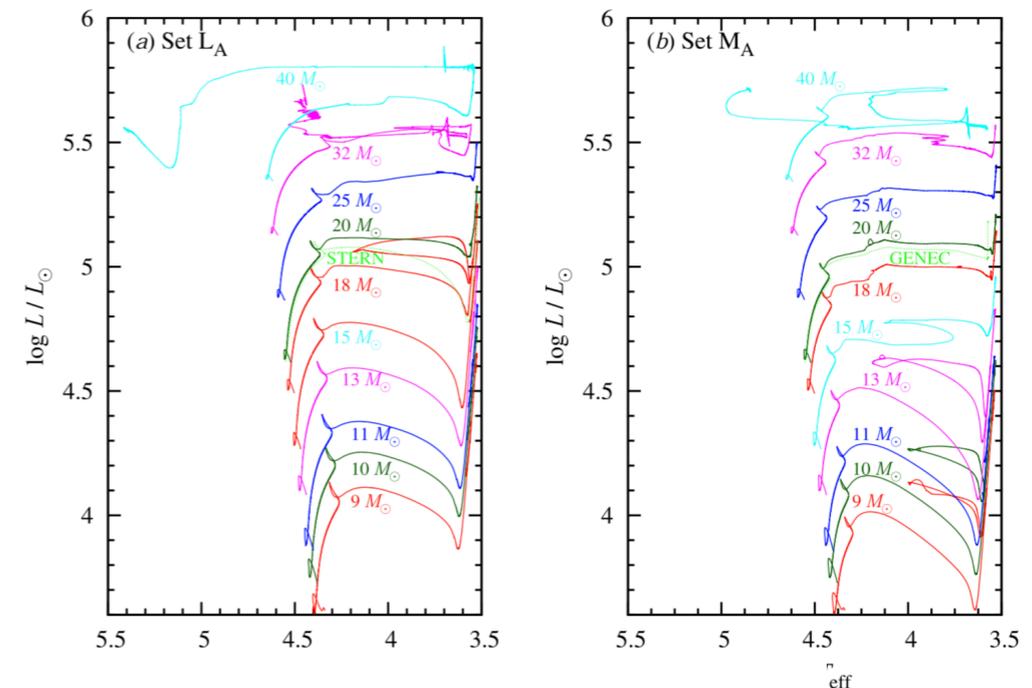
- Pop. III star with $M_{\text{zams}} = 85M_{\odot}$
 - Weak stellar wind mass loss
 - Expansion up to $\sim 160R_{\odot}$
 - He core with $\lesssim 40M_{\odot}$
- GW190521 can be formed from a Pop. III binary !!!



Uncertainty of Pop. III model

Yoshida et al. (2019)

- No massive Pop. III star is discovered so far.
- Extrapolation from nearby stars to Pop. III stars
- Nearby star models
 - AB-type stars in MW open clusters, **M model**
 - Early B-type stars in LMC, Stern (Brott et al. 2011) **L model**
- The maximum radius of a $80M_{\odot}$ star
 - M model: $\sim 40R_{\odot}$, similar to Farrell et al. (2020)
 - L model: $\sim 3 \times 10^3 R_{\odot}$, similar to Yoon et al. (2012)
- If the L model is correct, a Pop. III binary cannot form GW190521, the same as Pop. I/II binaries.

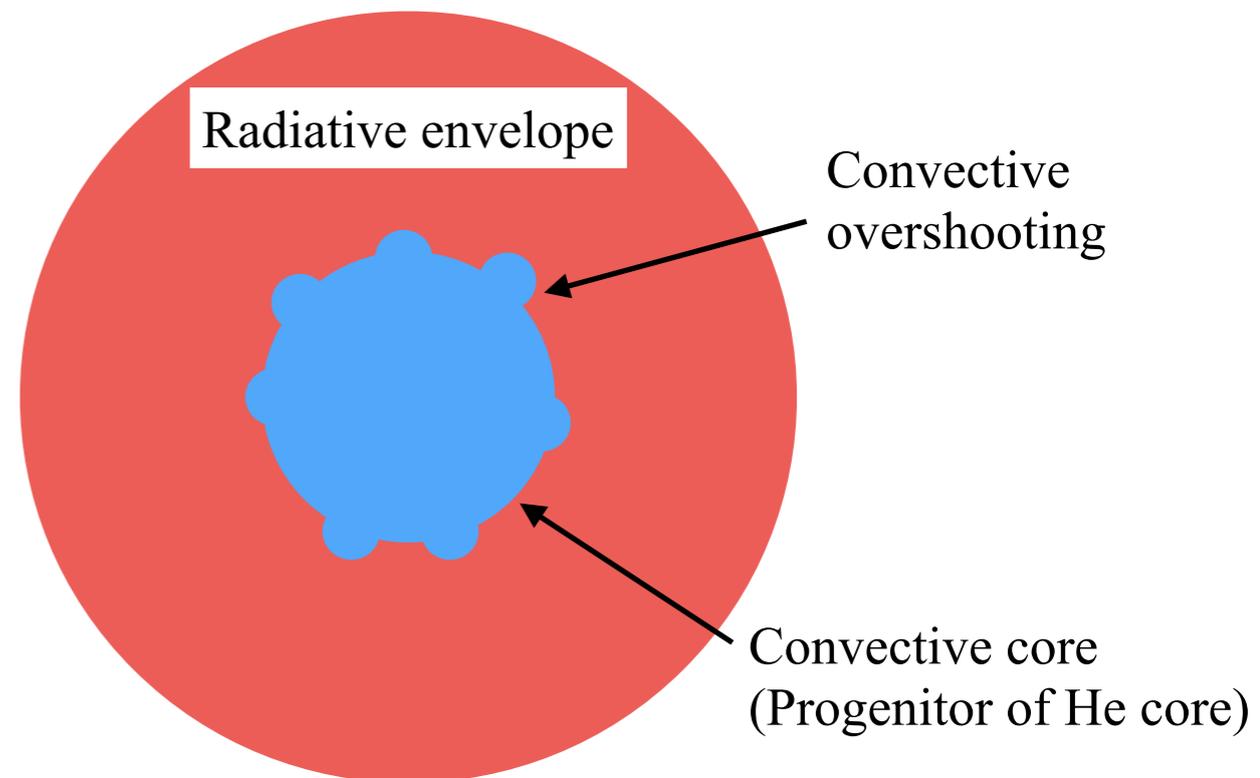
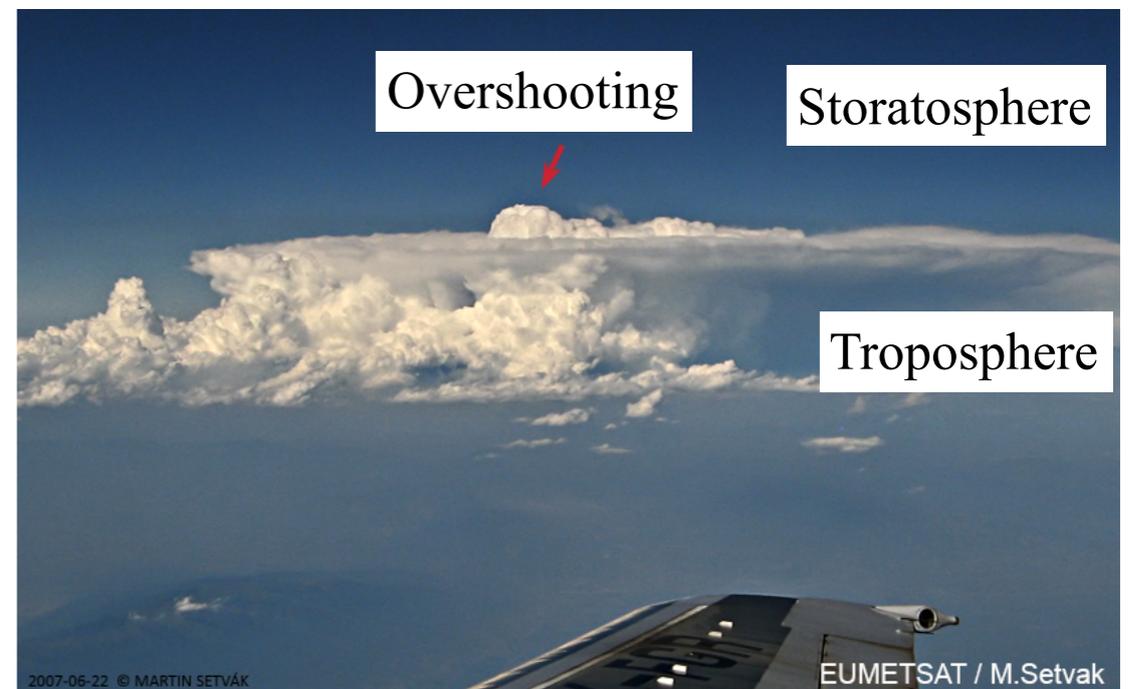


Two Pop. III models

Tanikawa et al. (2020c)

Convective overshooting

- Overshoot parameter: $f_{\text{ov}} \sim 0.02$
(Kippenhahn et al. 1990; 2012)
 - $D(z) = D_0 \exp \frac{-2z}{f_{\text{ov}} H_P}$
 - M model: $f_{\text{ov}} = 0.01$
 - L model: $f_{\text{ov}} = 0.03$
- Larger overshoot parameter (more effective overshooting)
 - Larger He core at the end of MS
 - Larger luminosity in post-MS
 - Larger radius in post-MS

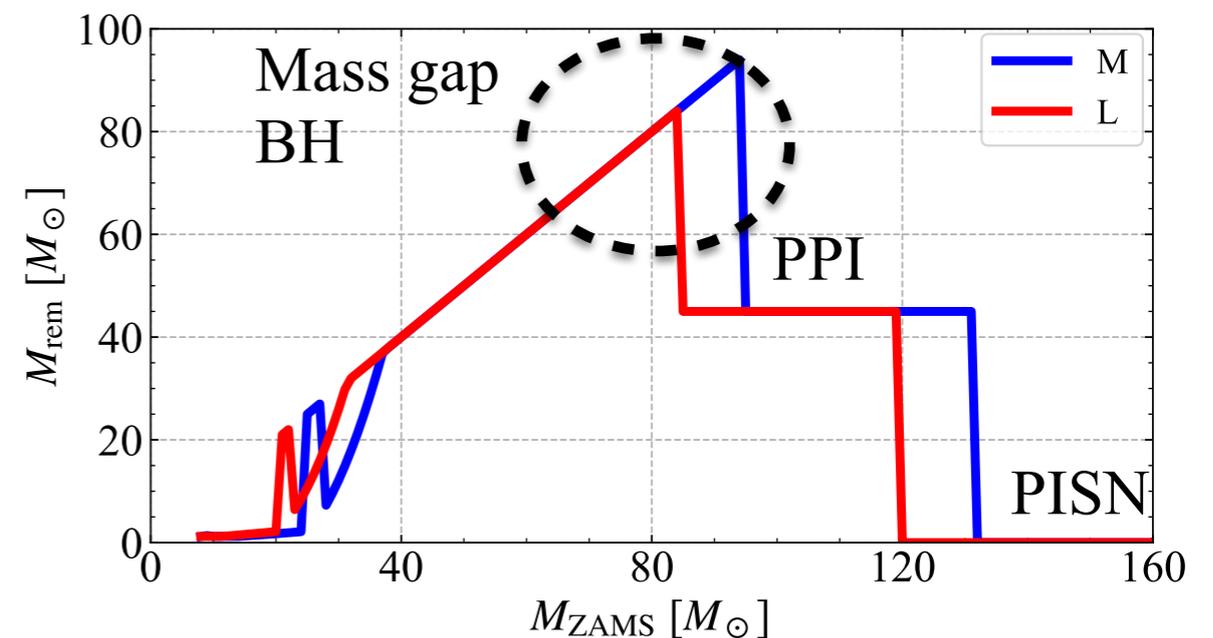
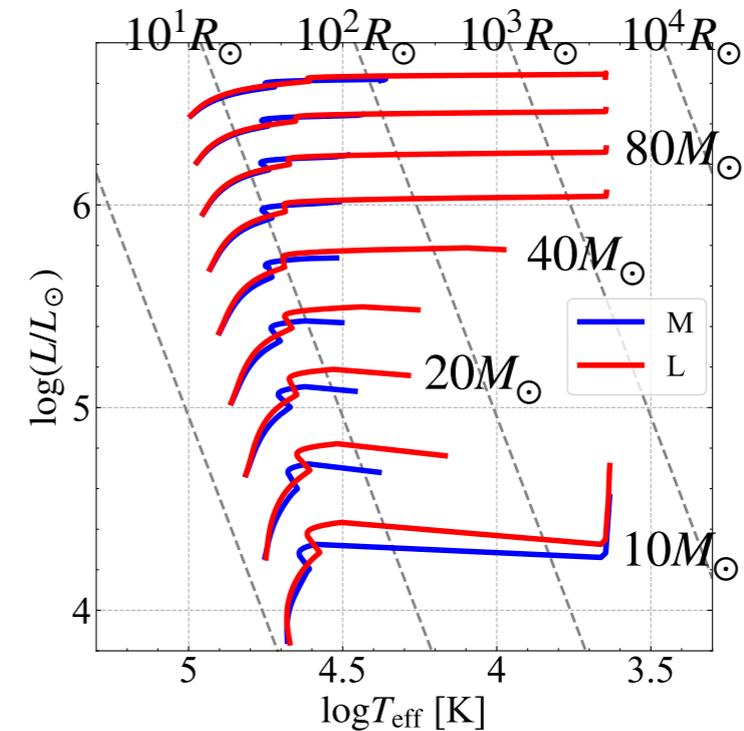


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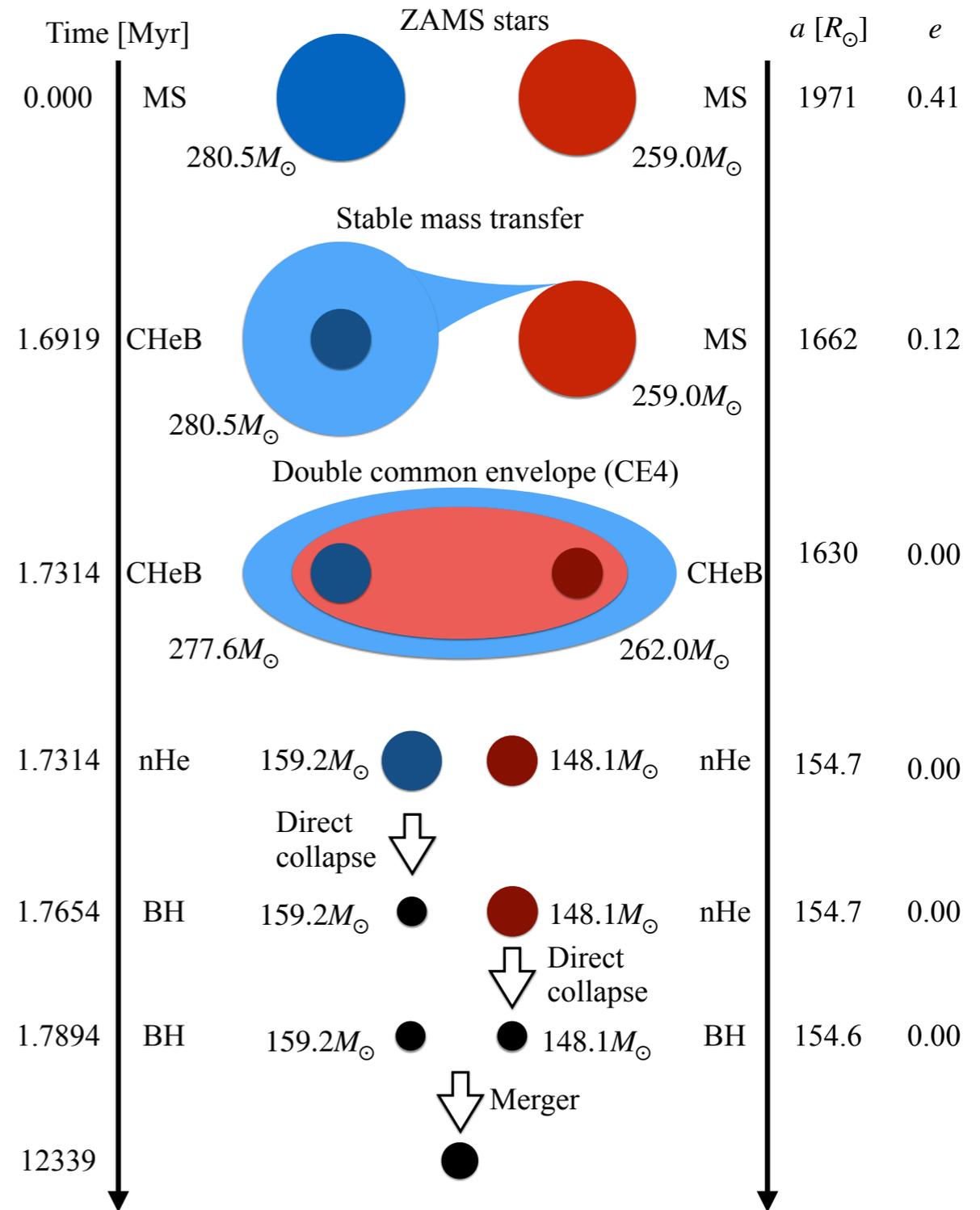
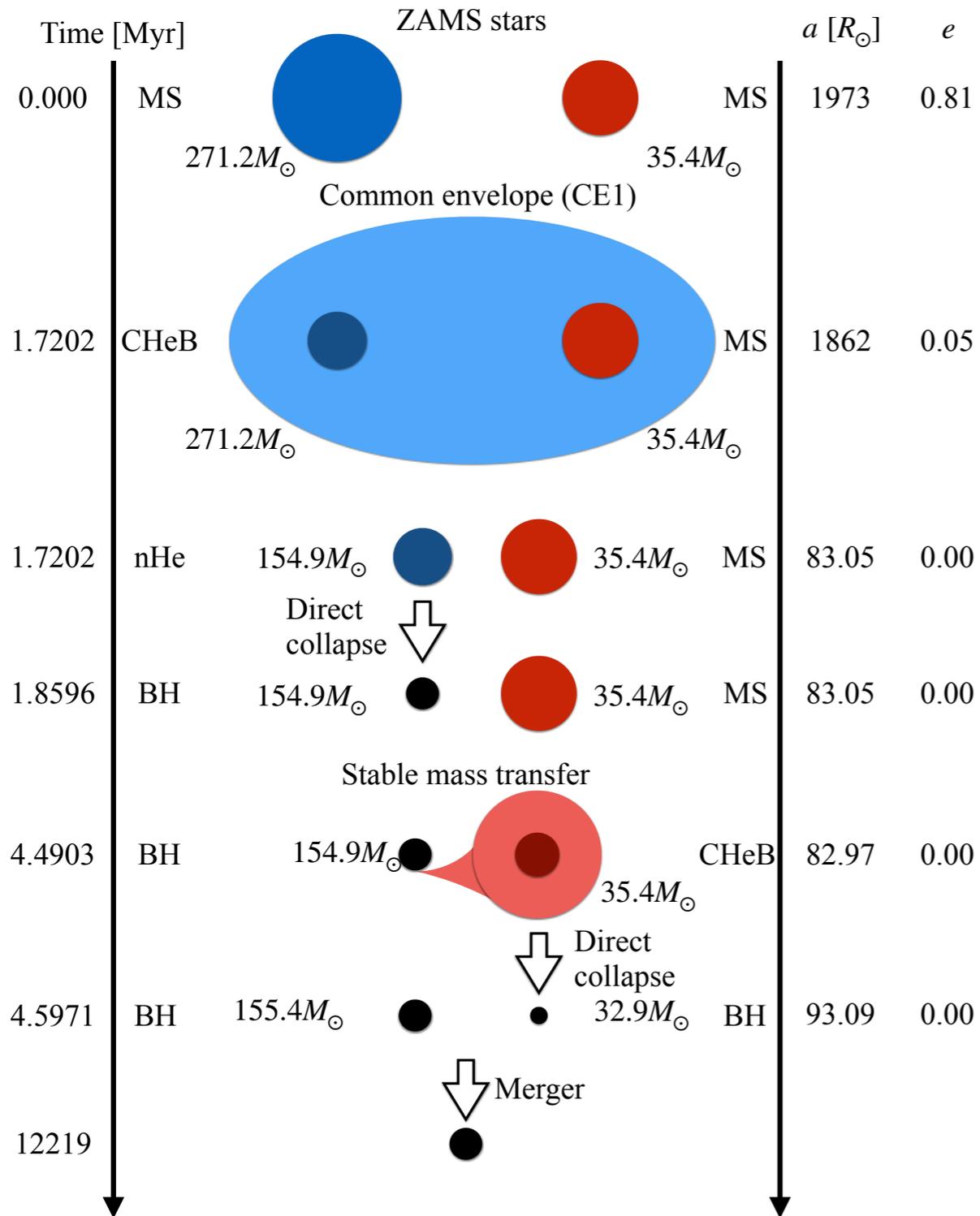
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Binary population synthesis

- BSE (Hurley et al. 2000; 2002) modified by Tanikawa et al. (2020a)
- Single star evolution
 - Fryer's rapid model with PPI/PISN
 - No stellar wind nor BH natal kick
- Binary star evolution
 - Tidal interaction
 - Stable mass transfer, common envelope
 - GW orbital decay
 - Etc.
- Initial conditions
 - $f(m_1) \propto m_1^{-1}, f(q) \propto \text{const}, f(a) \propto a^{-1}, f(e) \propto e$
- Cumulative Pop. III density
 - $\sim 10^{13} M_\odot \text{pc}^{-3}$ comparable to Magg et al. (2016) and Skinner, Wise (2020)

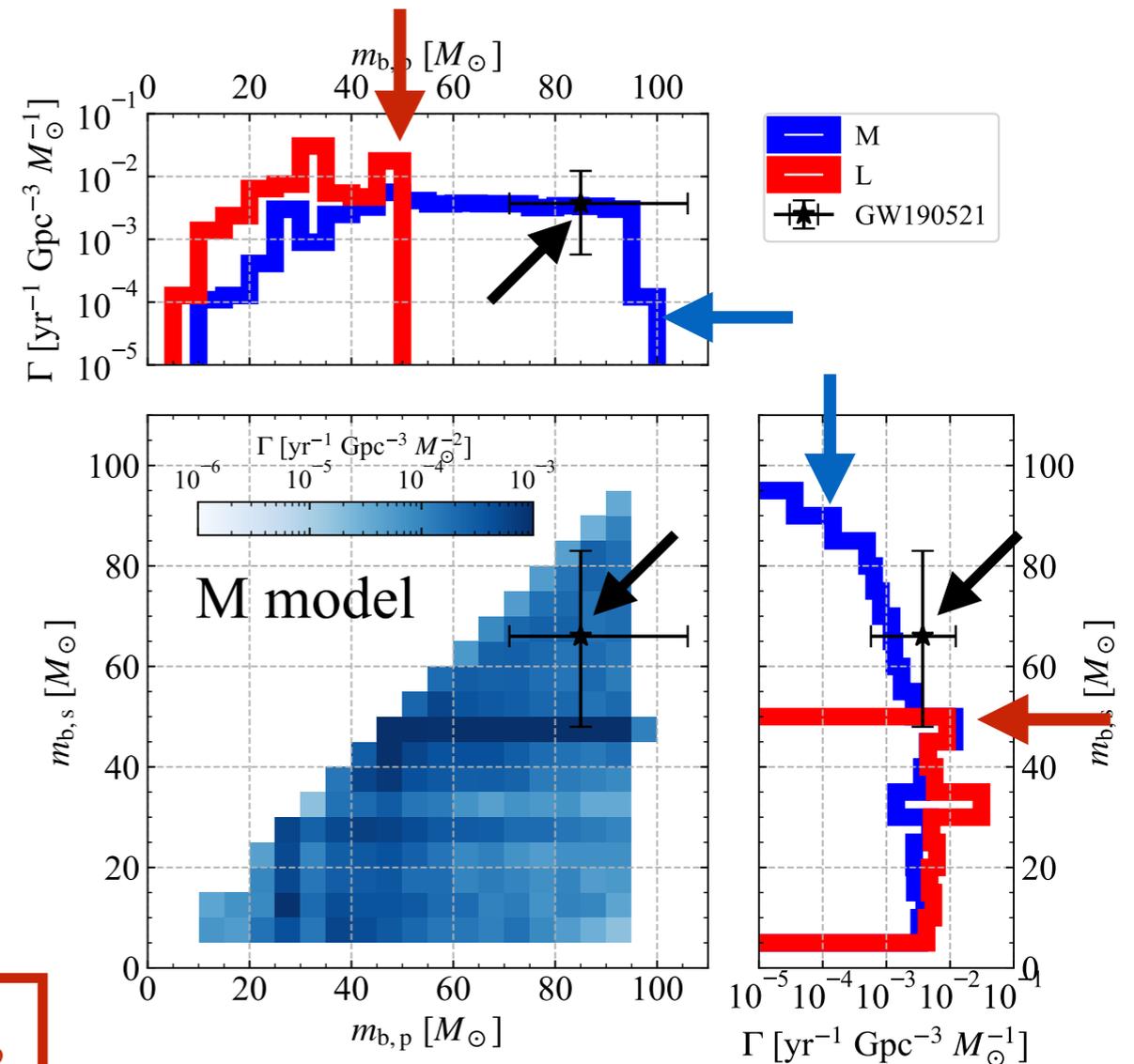


Introduction of BPS



BH mass distribution

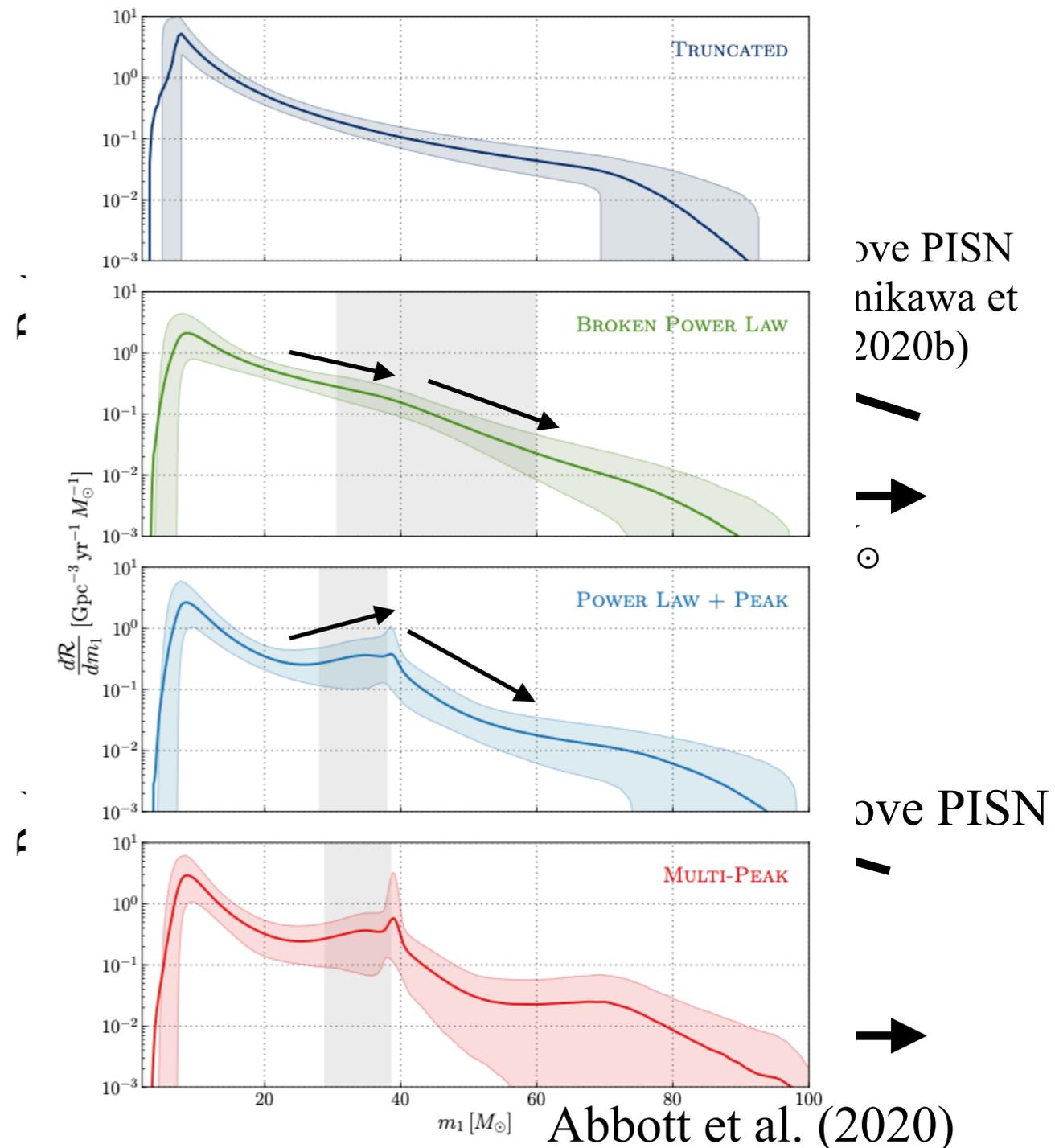
- M model
 - The maximum mass: $\sim 100M_{\odot}$
 - Stars lose little mass through binary interactions.
 - Pop. III stars can form GW190521-like BH-BHs.
 - Support for the claims of Farrell et al. (2020) and Kinugawa et al. (2020)
- L model
 - The maximum mass: $\sim 50M_{\odot}$
 - Stars lose their H envelopes through binary interactions
 - No Pop. III stars can form GW190521-like BH-BHs.



It depends on the choice of Pop. III models, or overshoot parameters

If GW190521 is Pop.III ...

- Even if the M model is correct, no Pop. III binary can form BH-BHs with $100 - 130M_{\odot}$.
- If GW190521 is Pop. III, the merger rate of BH-BHs with $100 - 130M_{\odot}$ is much smaller than with $50 - 100M_{\odot}$.
- But, the converse is not true.
 - A Pop. II binary can form GW190521 if the $^{12}\text{C}(\alpha, \gamma)^{16}\text{O}$ reaction rate is 3 times smaller than the standard rate (Belczynski 2020).

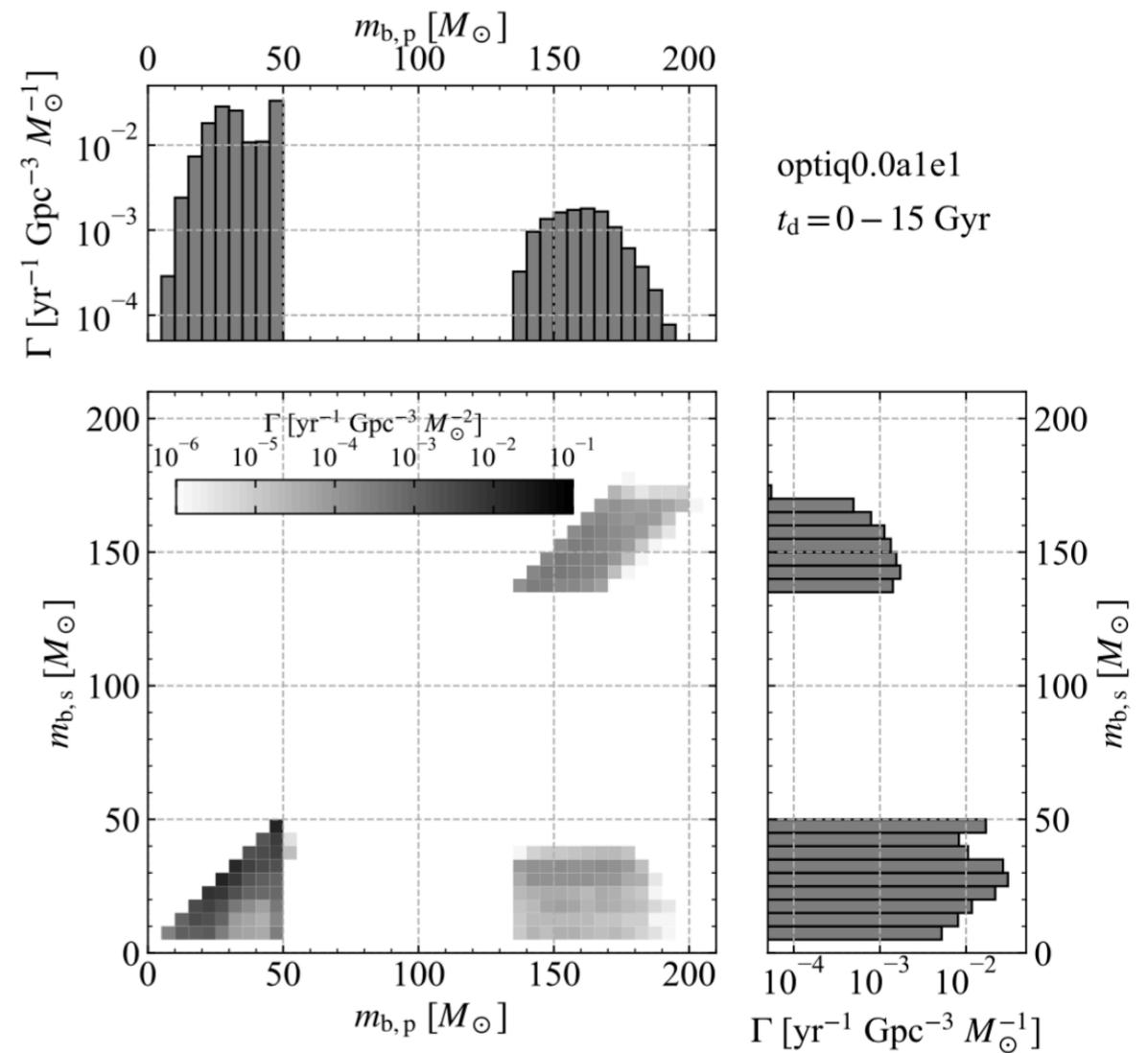


Summary

- GW190521 is a merger of a BH-BH with at least a BH in the PI mass gap.
- GW190521 can be formed from a Pop. III binary (Farrell et al. 2020; Kinugawa et al. 2020)
- We have shown that the Pop. III scenario strongly depends on the effectiveness of convective overshooting.
- If GW190521 is a Pop. III origin, the merger rate of $100 - 130M_{\odot}$ BHs is much smaller than that of $50 - 100M_{\odot}$.
 - But, the converse is not true.

Lモデル詳細

- Initial conditions
 - $m_{1, \max} = 150M_{\odot} \rightarrow 300M_{\odot}$
- $130M_{\odot}$ 以上のBH形成
- BH Mergers of $> 130M_{\odot}$ and $< 50M_{\odot}$ are $\sim 10^{-2} \text{Gpc}^{-3} \text{yr}^{-1}$.
- GW190521 could be below and above the mass gap (Fishbach, Holtz 2020; Nitz, Capano 2020).



Tanikawa et al. (2020b)