



連星ブラックホール形成：孤立連星と球状星団

Binary black hole formation:
binary evolution and dynamical capture

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初代星・初代銀河研究会, First star and first galaxy

2020年11月, 東北大学, Nov. 2020, Tohoku U.

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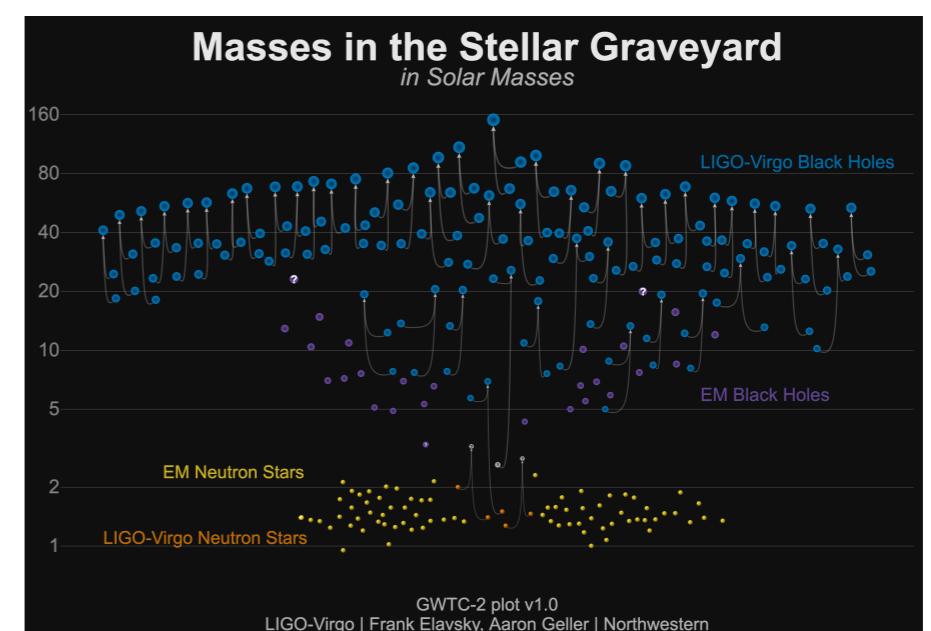
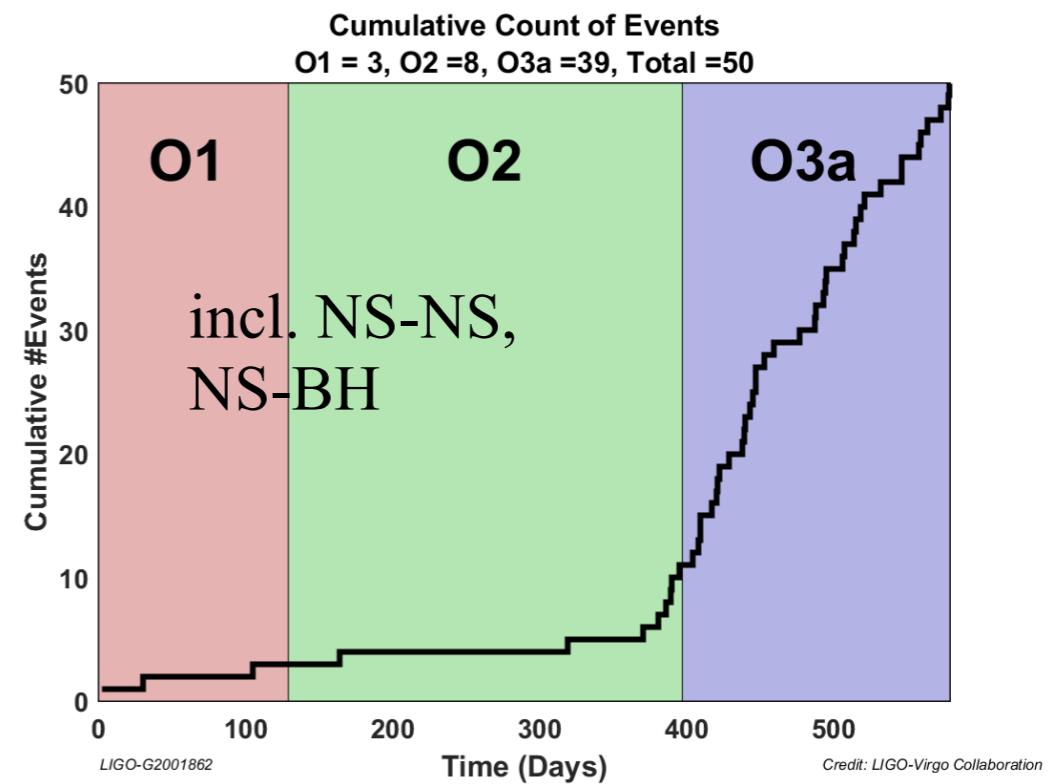
- From O1/O2 to O3a
- Globular clusters
- Pop. III binaries

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- Globular clusters
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GW events

- The first detection GW150914 in 2015 (only LIGO)
- 10 BH-BHs discovered in O1/O2 by 2017 (entry of Virgo)
- 44 confident BH-BHs in O1/O2/ O3a by the last month
- GW sensitivity to NS-NSs
 - H: 66Mpc to 108Mpc
 - L: 88Mpc to 135Mpc
 - V: 26Mpc to 45Mpc

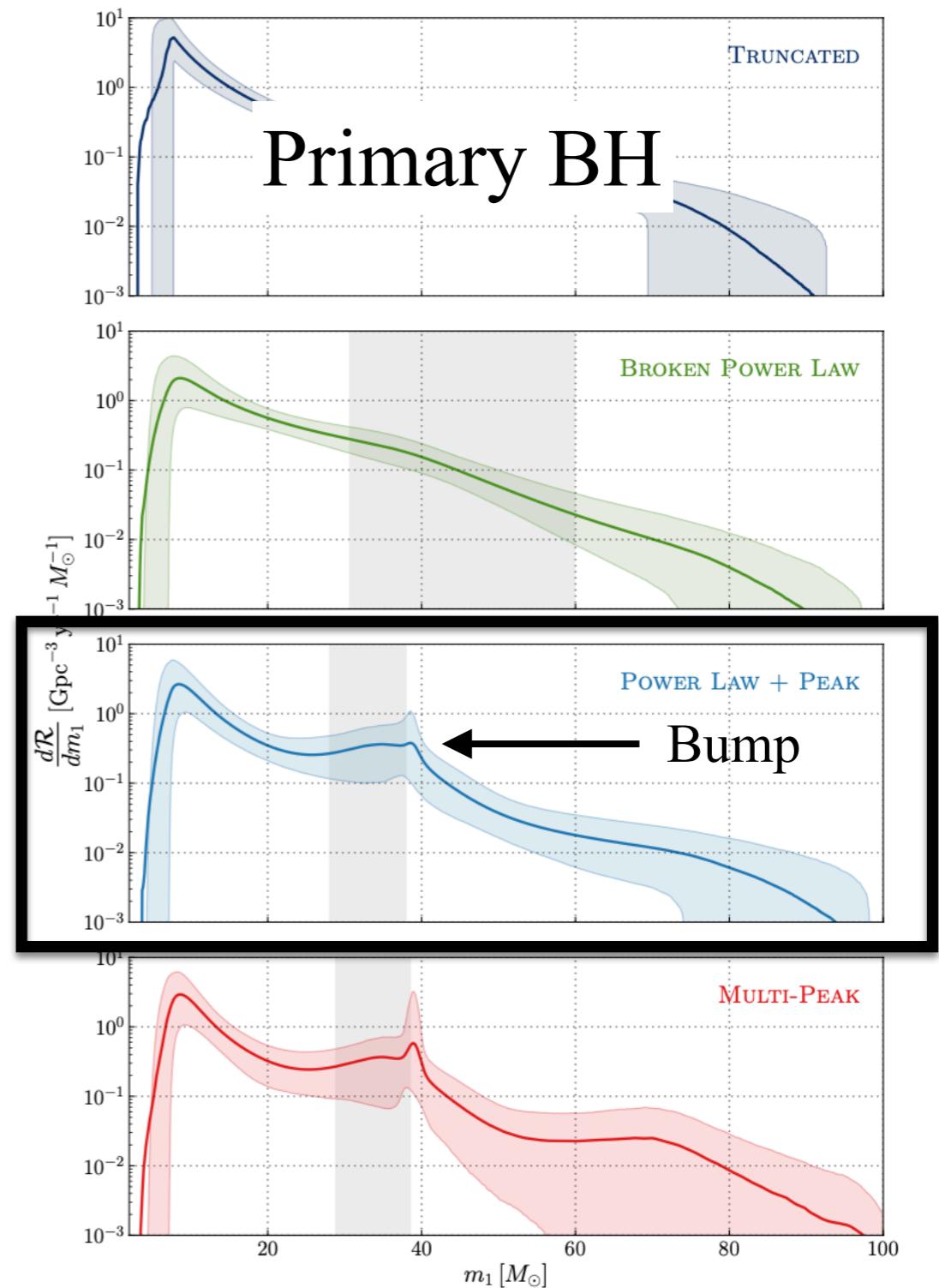


BH-BH rate density

- O1/O2: $53.2^{+55.8}_{-28.2} \text{ Gpc}^{-3} \text{ yr}^{-1}$
- O1/O2/O3a
 - No cosmic evolution: $23.9^{+14.9}_{-8.6} \text{ Gpc}^{-3} \text{ yr}^{-1}$
 - Cosmic evolution: $19.7^{+57.3}_{-15.9} \text{ Gpc}^{-3} \text{ yr}^{-1}$ at $z = 0$
- $\sim 10 - 100 \text{ Gpc}^{-3} \text{ yr}^{-1} \rightarrow \sim 10 \text{ Gpc}^{-3} \text{ yr}^{-1}$
- Good news for scenarios other than binary evolution?

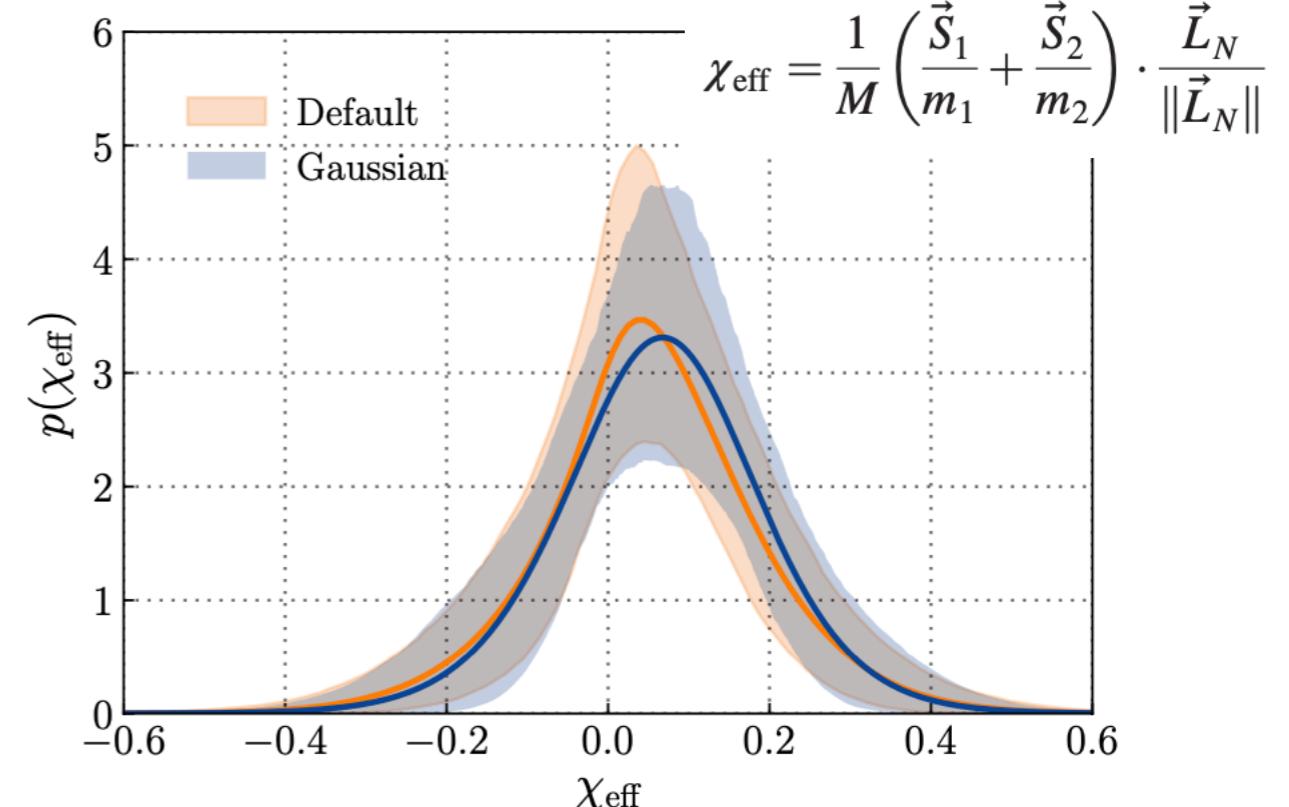
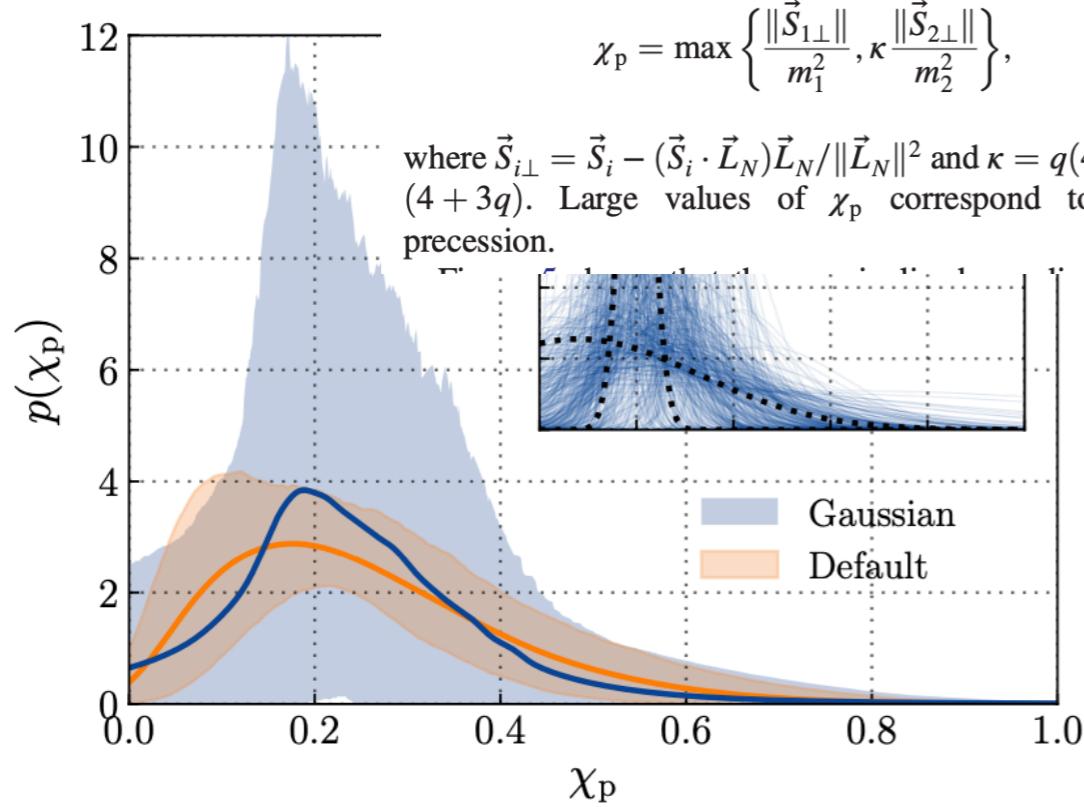
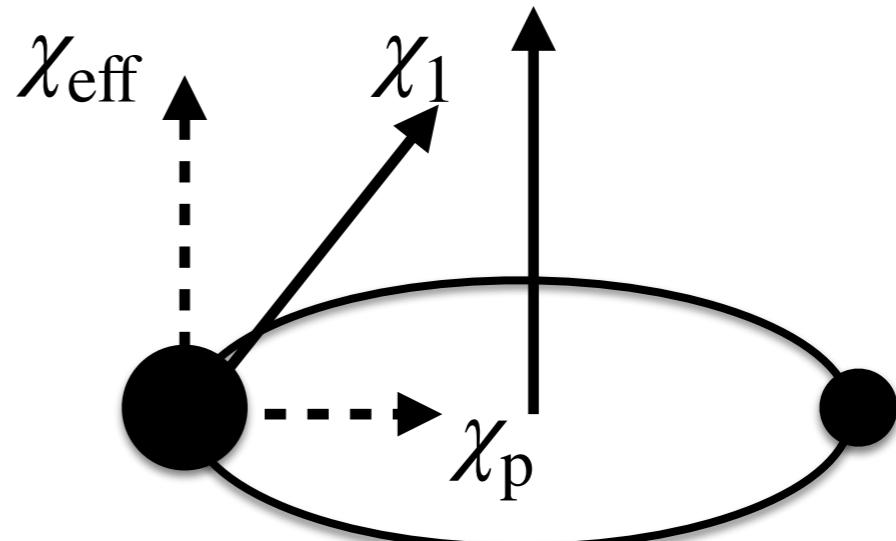
BH mass distribution

- Primary BH
 - Global maximum at $7.8^{+2.2}_{-2.1} M_{\odot}$
 - Lower mass gap inferred by BH X-ray binaries
 - Break or bump at $\sim 40 M_{\odot}$
 - Pair instability (or Pop. III?)
- Mass ratio
 - Consistent with $q = m_2/m_1 \sim 1$
 - But, see GW190412, GW190426_152155 (NS-BH?) and GW190814 (NS-BH?)



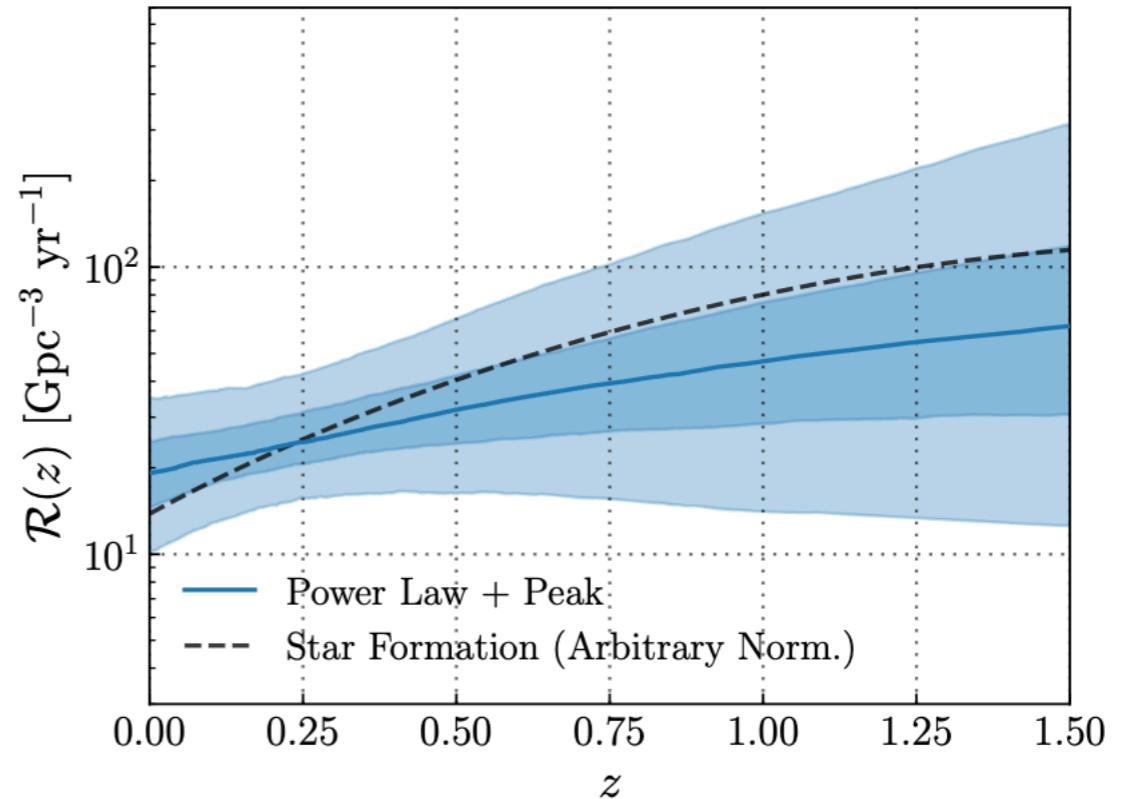
BH spin distribution

- 12-44% of BH-BH spins tilted by $> 90^\circ$
 - ~ 500 km/s BH kick?
 - Intrinsic spin?
 - Dynamical capture?
- No single event with $\chi_p > 0$ nor $\chi_{\text{eff}} < 0$



Cosmic evolution

- $\mathcal{R}_{\text{BH-BH}} \propto (1 + z)^\kappa$
- Consistent with a non-evolving distribution ($\kappa = 0$)
- But, more consistent with $\kappa = 1.3^{+2.1}_{-2.1}$ and $\kappa = 1.8^{+2.1}_{-2.2}$
- Slower than the star formation rate ($\kappa = 2.7$)
- BH-BHs with long delay time



BH-BH formation scenarios

- Binary evolution (Kinugawa's talk)
 - Pop. I/II common envelope evolution
 - Pop. I/II chemically homogeneous evolution
 - Pop. III stable mass transfer
- Dynamical capture in dense stellar clusters
 - Globular cluster (Alessandro's and Long's talks)
 - Open cluster (Kumamoto's and Alessandro's talks)
 - Galactic center (Tagawa's talk)
- Hierarchical multiple star evolution
- Primordial BH

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Black hole budget (GC)

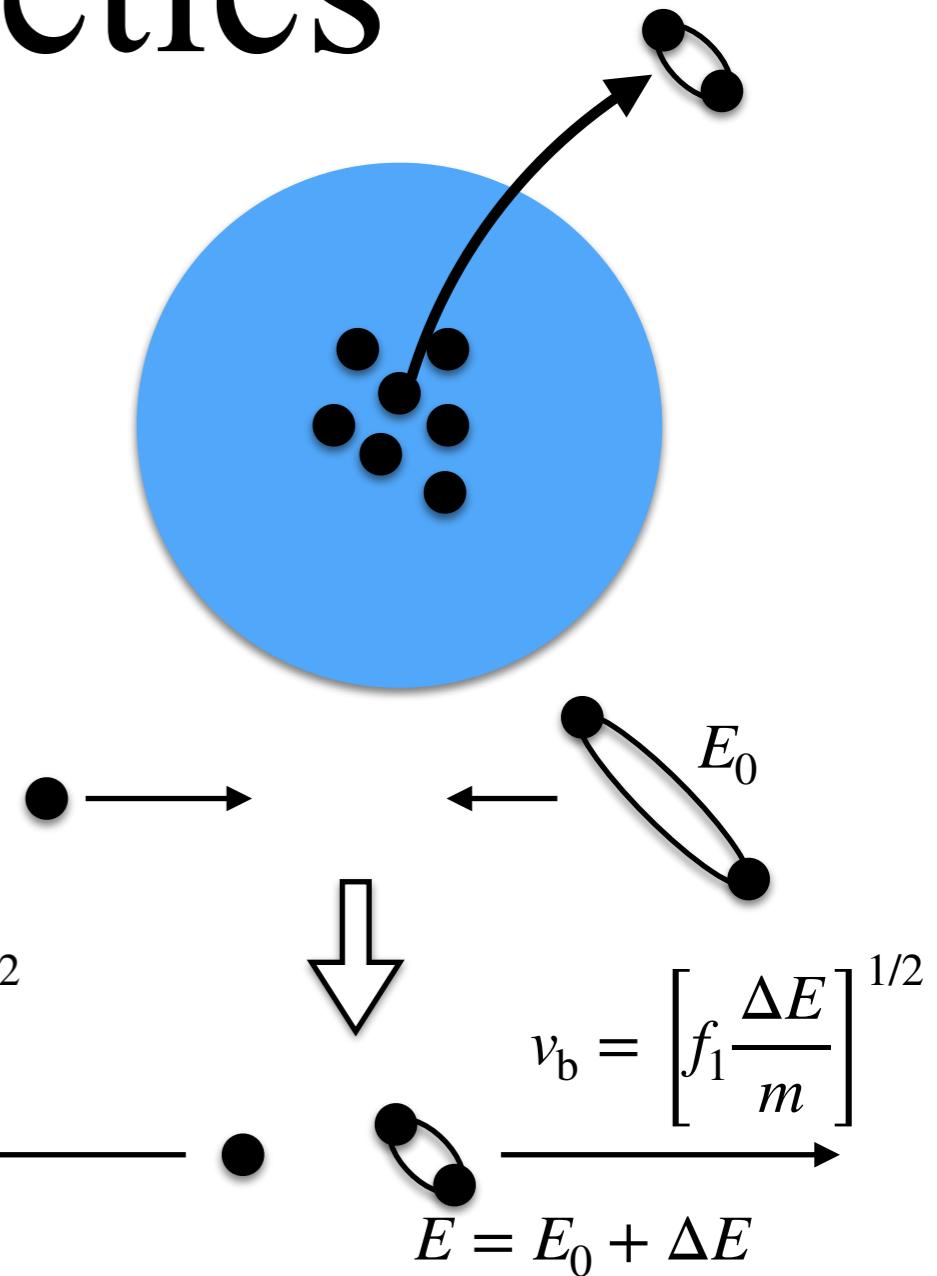
- BH-BH density: $n_{\text{BH-BH}} \sim 10^{11} \text{Gpc}^{-3} \left(\frac{\Gamma_{\text{BH-BH}}}{10 \text{Gpc}^{-3} \text{yr}^{-1}} \right) \left(\frac{T_{\text{Hubble}}}{10 \text{Gyr}} \right)$
 - Binary evolution:
 - $n_{\text{BH,star}} \sim 10^{15} \text{Gpc}^{-3} \left(\frac{\rho_{\text{star}}}{10^{18} M_{\odot} \text{Gpc}^{-3}} \right) \left(\frac{\eta_{\text{BH}}}{10^{-3} M_{\odot}} \right)$
 - $\frac{n_{\text{BH-BH}}}{n_{\text{BH,star}}} \sim 10^{-4} \left(\frac{\Gamma_{\text{BH-BH}}}{10 \text{Gpc}^{-3} \text{yr}^{-1}} \right)$
 - Dynamical capture in globular clusters:
 - $n_{\text{BH,GC}} \sim 10^{12} \text{Gpc}^{-3} \left(\frac{\rho_{\text{star}}}{10^{18} M_{\odot} \text{Gpc}^{-3}} \right) \left(\frac{\eta_{\text{BH}}}{10^{-3} M_{\odot}} \right) \left(\frac{\rho_{\text{GC}}/\rho_{\text{star}}}{10^{-3}} \right)$
 - $\frac{n_{\text{BH-BH}}}{n_{\text{BH,GC}}} \sim 0.1 \left(\frac{\Gamma_{\text{BH-BH}}}{10 \text{Gpc}^{-3} \text{yr}^{-1}} \right) \left(\frac{\rho_{\text{GC}}/\rho_{\text{star}}}{10^{-3}} \right)$
- Abbott et al. (2020)
- Kroupa (2001), $M_{\text{zams}} \gtrsim 20 M_{\odot} \rightarrow \text{BH}$
- Madau, Dickinson (2014)
- MW GC

Supplementary budget

- More many GCs at the formation time ($\rho_{\text{GC}}/\rho_{\text{star}} > 10^{-3}$)
 - Gas expulsion
 - Stellar evolution mass loss (stellar wind, supernova)
 - Thermodynamical evaporation
- Repeated mergers (e.g. Rodriguez et al. 2019)
- Dark clusters (e.g. Wang 2020)

BH-BH Energetics

- $t_{\text{gw}} = \frac{5}{256} \frac{c^5}{G^3} \frac{a^4}{m_1 m_2 (m_1 + m_2)} g(e)$
 - $g(e) = \frac{(1 - e^2)^{3.5}}{1 + (73/24)e^2 + (37/96)e^4}$
- $a_{\text{esc}} = \frac{f_1 f_2}{2} \frac{Gm}{v_{\text{esc}}^2}$
 - $v^2 = f_1 \frac{\Delta E}{m}$
 - $\Delta E = f_2 E = f_2 \frac{Gm^2}{2a}$
- $t_{\text{gw}} = 12 \text{Gyr} \left(\frac{f_1}{1/3} \right)^4 \left(\frac{f_2}{1/1.4} \right)^4 \left(\frac{m}{30M_\odot} \right) \left(\frac{v_{\text{esc}}}{50 \text{km/s}} \right)^{-8} \left(\frac{g(e)}{g(0.9)} \right)$

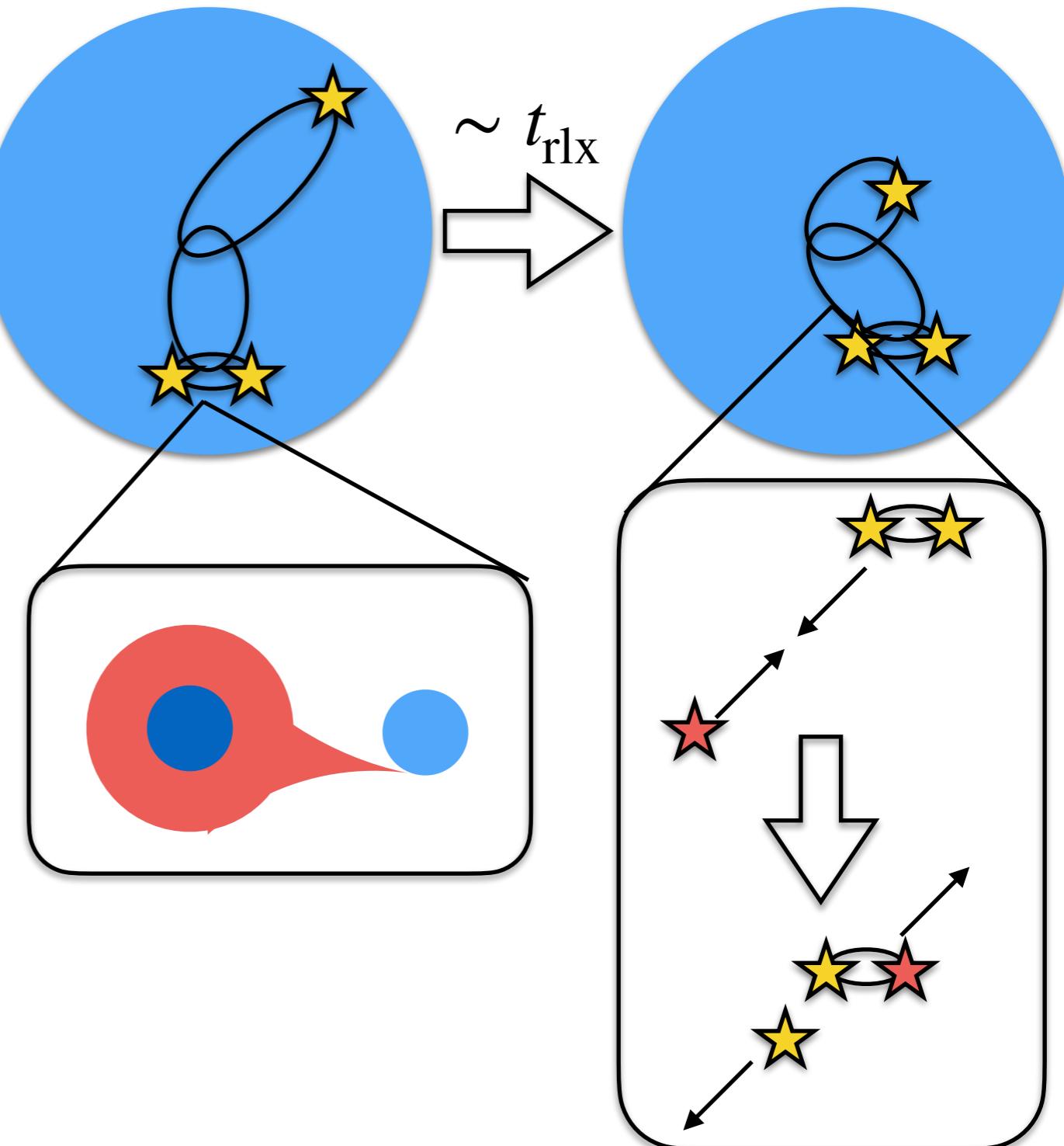


Numerical simulation

- N-body simulation (Portegies Zwart, McMillan 2000; Banerjee et al. 2010; Tanikawa 2013; Fujii et al. 2017; etc)
 - Small N ($\lesssim 10^5$), but $N \sim 10^6$ in reality
- Toy model (O’Leary et al. 2006; Sadowski et al. 2008; etc)
 - No cluster evolution
- Monte Carlo method (Downing et al. 2010; Rodriguez et al. 2016; Askar et al. 2017; etc)
 - $N \sim 10^6$, but approximate evolution

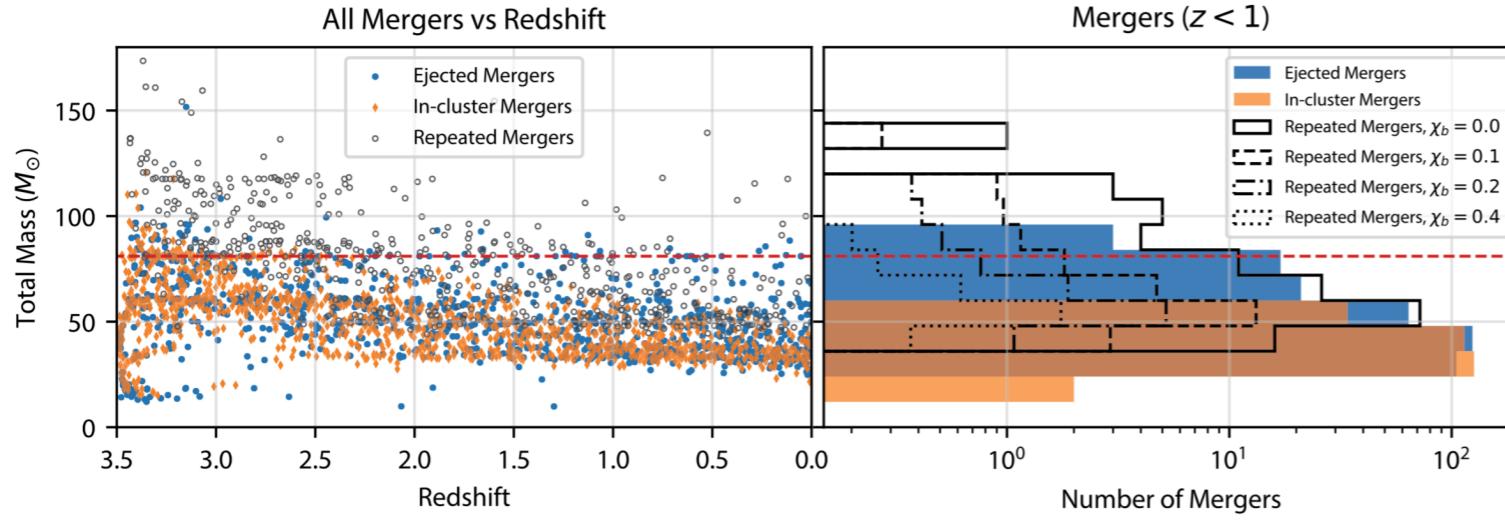
Monte Carlo method

- Single and binary stars orbit in a GC with constant E and J during two-body relaxation time.
- Their E and J diffuse every two-body relaxation time.
- They probabilistically experience close encounters with other single or binary stars, which calculated as gravitational few-body problems.
- They evolve due to stellar evolution and binary evolution, such as tidal interaction, mass transfer, common envelope evolution, and so on.
- Two codes: CMC (Rodriguez), MOCCA (Giersz)
- Several problems are becoming clear (Long 's talk).

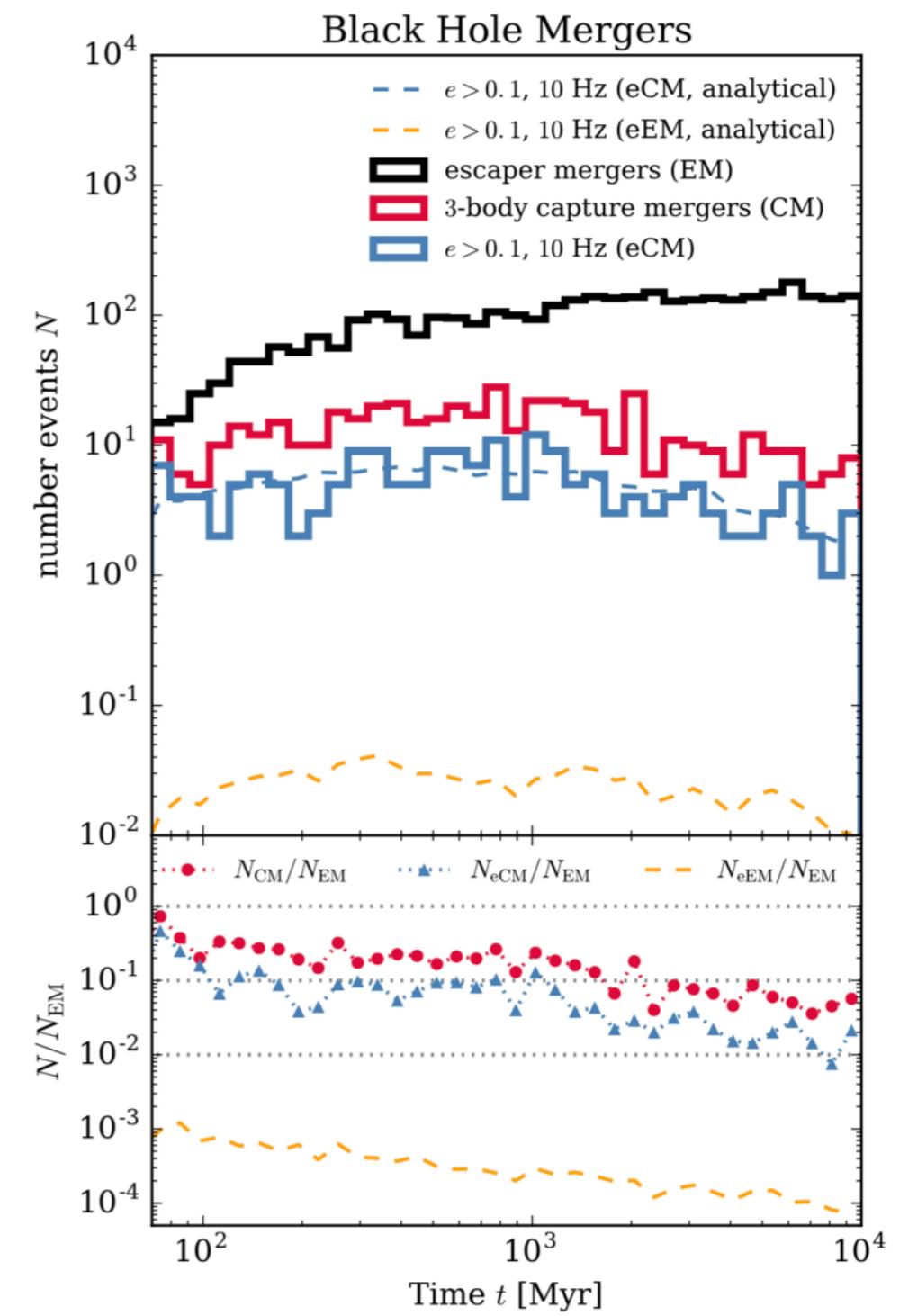


In-cluster mergers

- 10% of BH-BHs merge inside of GCs owing to high eccentricity.
 - The MOCCA code (right)
 - The CMC code (below)
- This allows repeated mergers.
 - BHs with large masses and spins
 - Filling of pair-instability mass gap



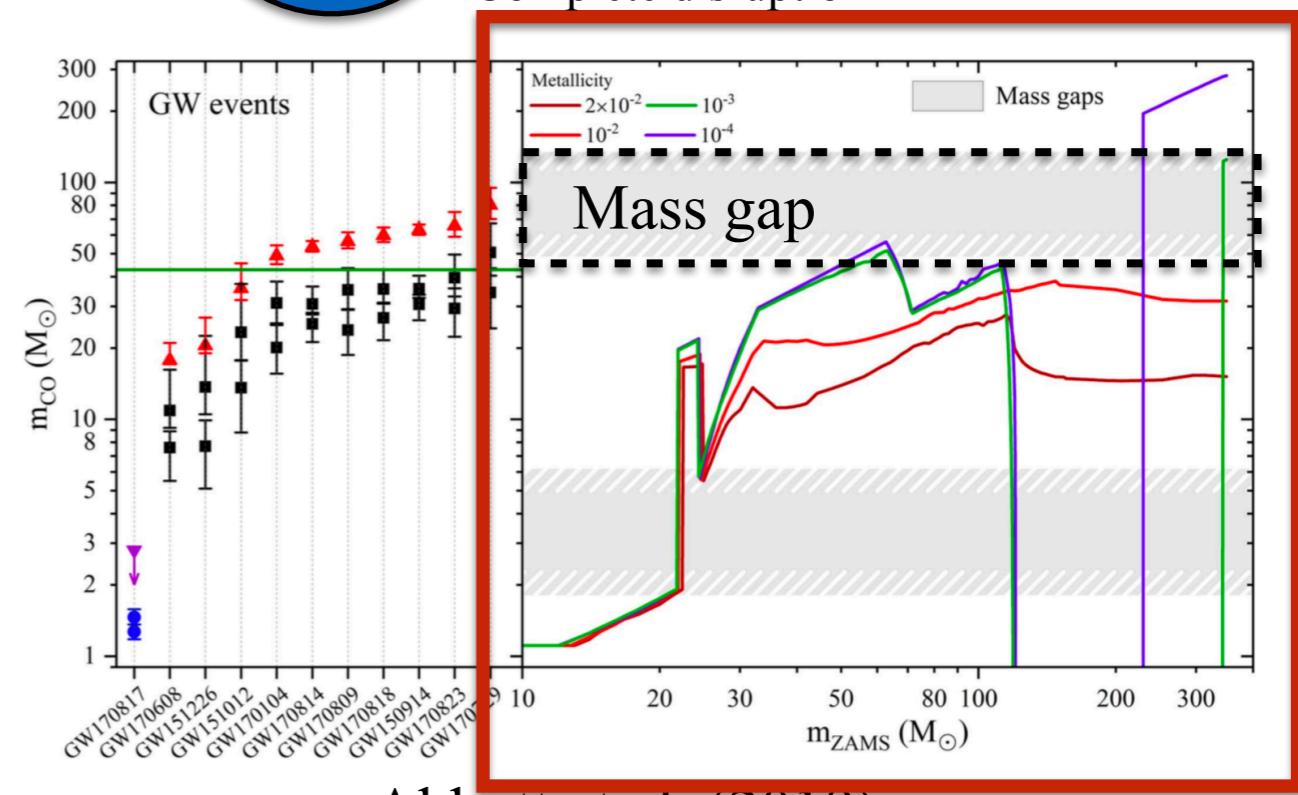
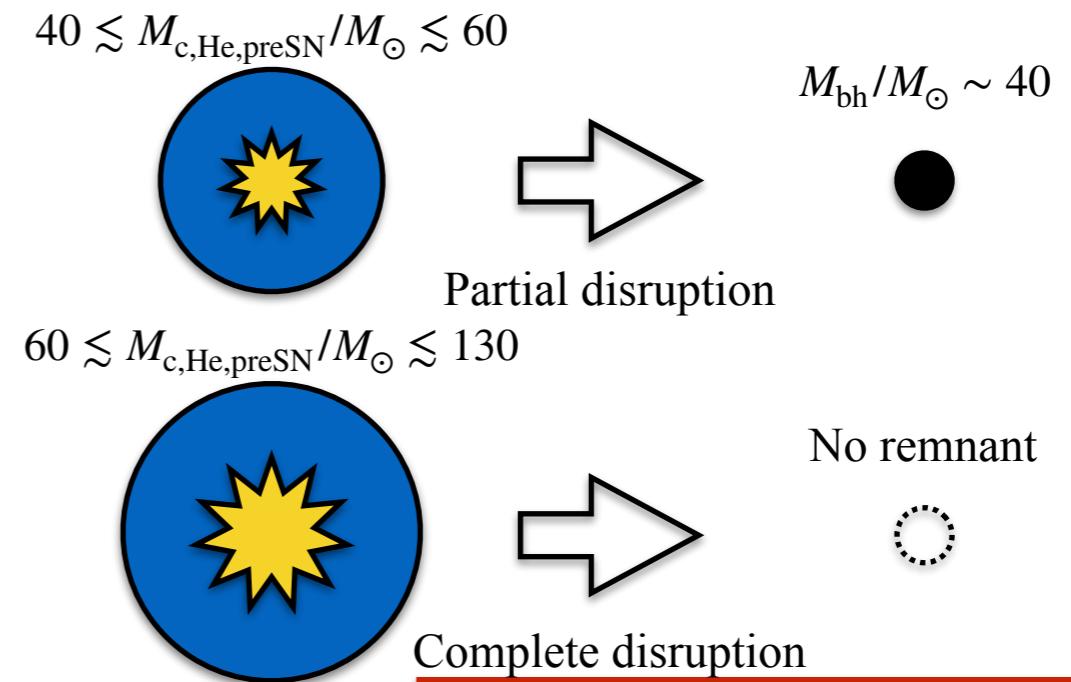
Rodriguez et al. (2018)



Samsing et al. (2018)

PPI and PISN

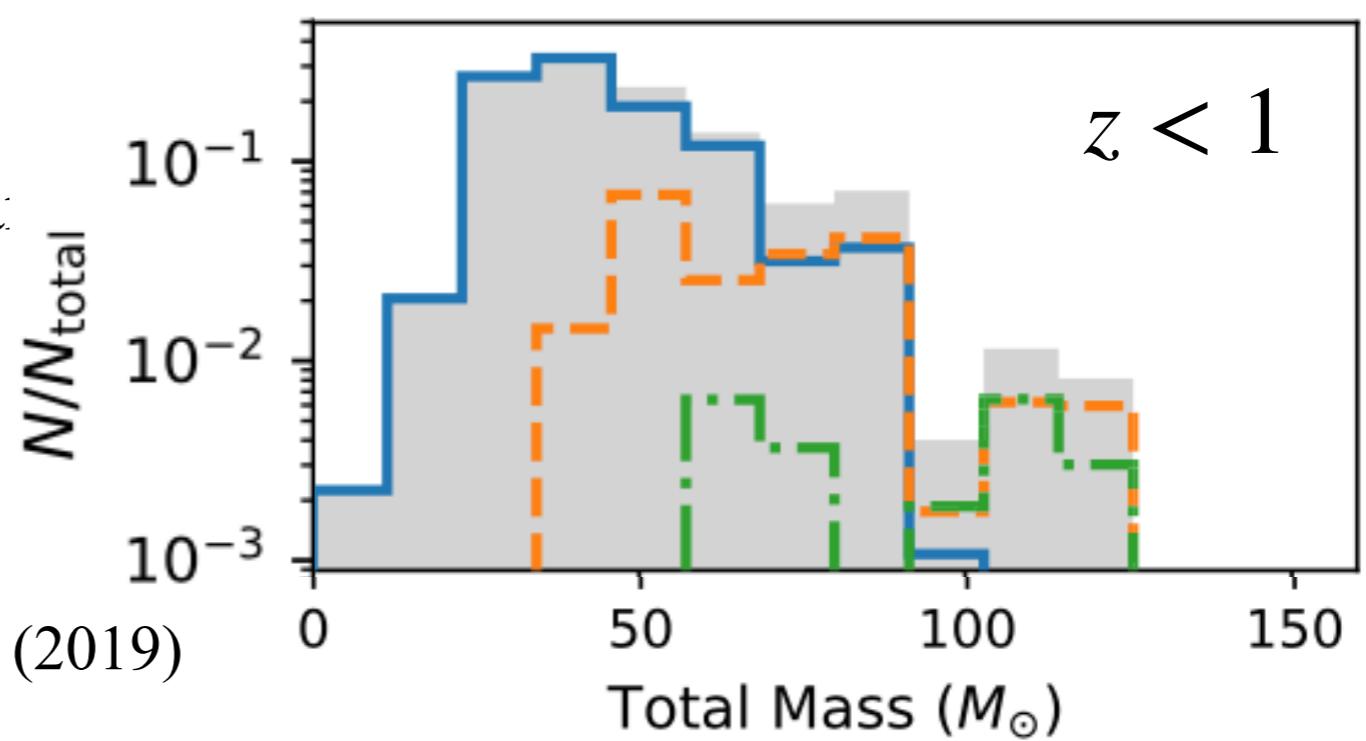
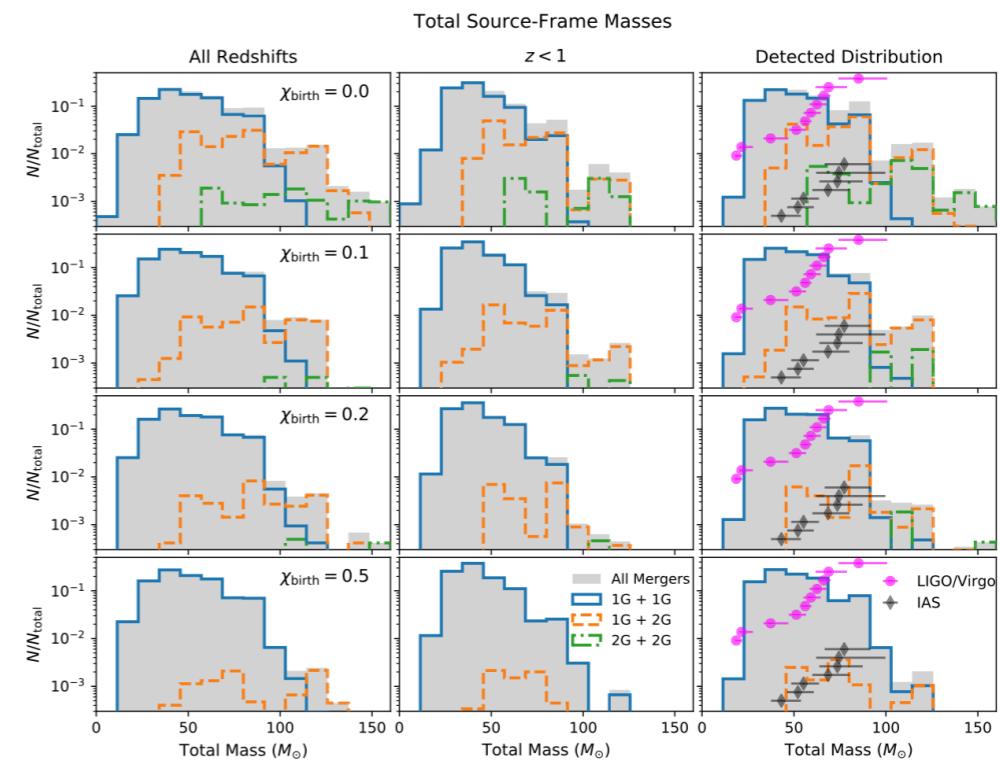
- Pulsational Pair Instability (PPI)
 - $40 \lesssim M_{\text{c,He,preSN}}/M_{\odot} \lesssim 60$
 - He core partially disrupted
 - $M_{\text{bh}} \sim 40M_{\odot}$
- Pair instability supernova (PISN)
 - $60 \lesssim M_{\text{c,He,preSN}}/M_{\odot} \lesssim 130$
 - He core completely disrupted
 - No remnant



Abbott et al. (2019)

BH Mass distribution

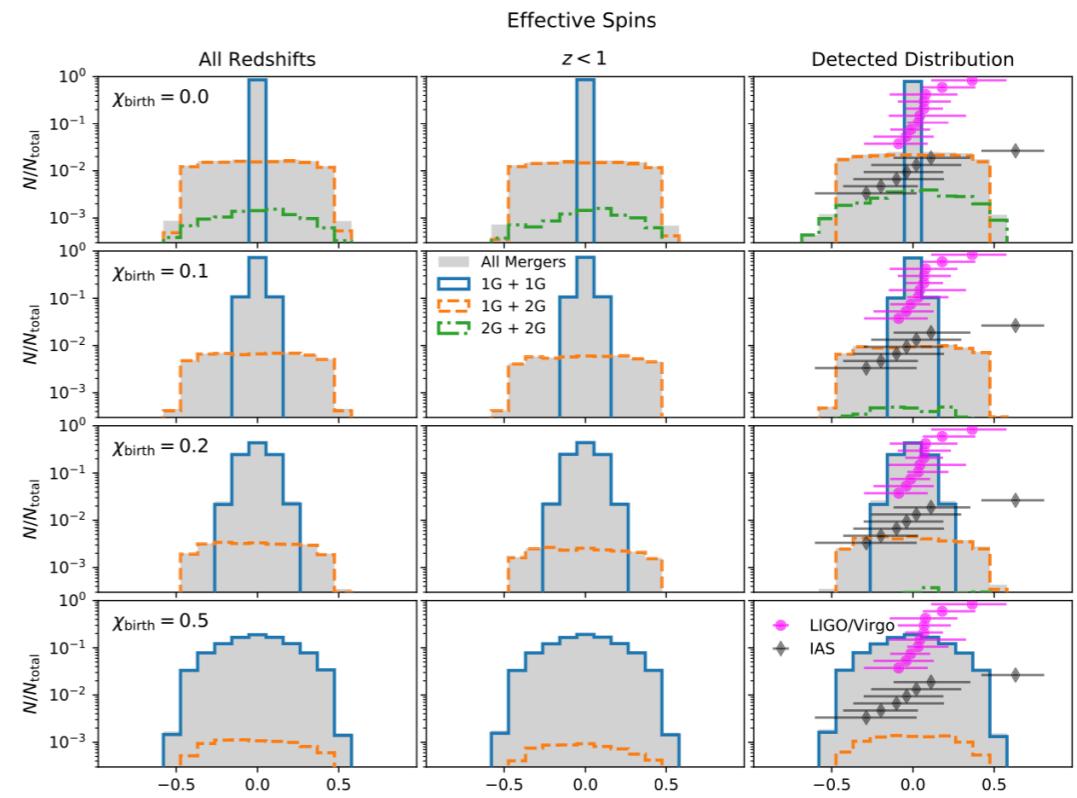
- $\frac{\mathcal{R}_{\text{BH-BH}}(M_{\text{tot}} = 150M_{\odot})}{\mathcal{R}_{\text{BH-BH}}(M_{\text{tot}} = 70M_{\odot})} \sim 10^{-1} - 10^{-1.5}$
 $\mathcal{R}_{\text{BH-BH}}(M_{\text{tot}} = 70M_{\odot})$
 at $z < 1$ inferred by GW observations
- $\frac{\mathcal{R}_{\text{BH-BH}}(M_{\text{tot}} = 150M_{\odot})}{\mathcal{R}_{\text{BH-BH}}(M_{\text{tot}} = 70M_{\odot})} < 10^{-2}$
 $\mathcal{R}_{\text{BH-BH}}(M_{\text{tot}} = 70M_{\odot})$
 at $z < 1$ for GCs
- Strong PPI/PISN model (No $> 40M_{\odot}$ BH without stellar a BH mergers)



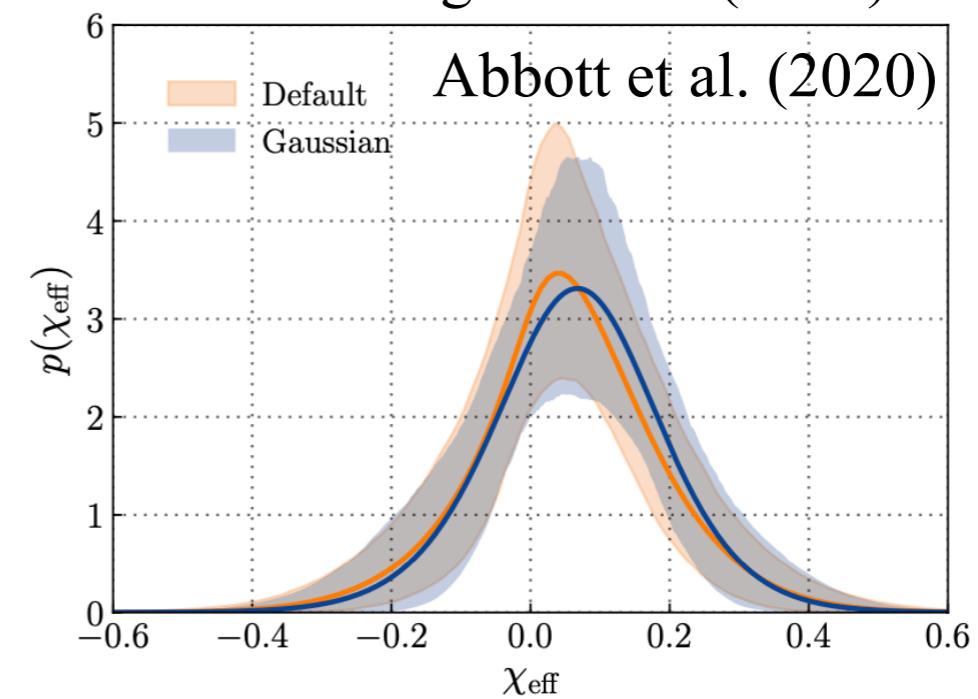
Rodriguez et al. (2019)

BH spin distribution

- Isotropic spin distribution for GCs
- GC BH-BHs may need other populations if they are the major component.

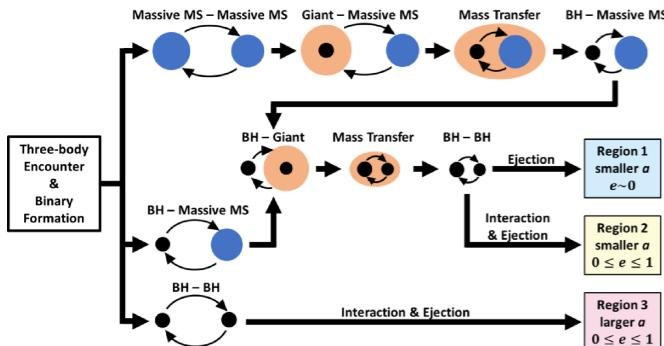


Rodriguez et al. (2019)

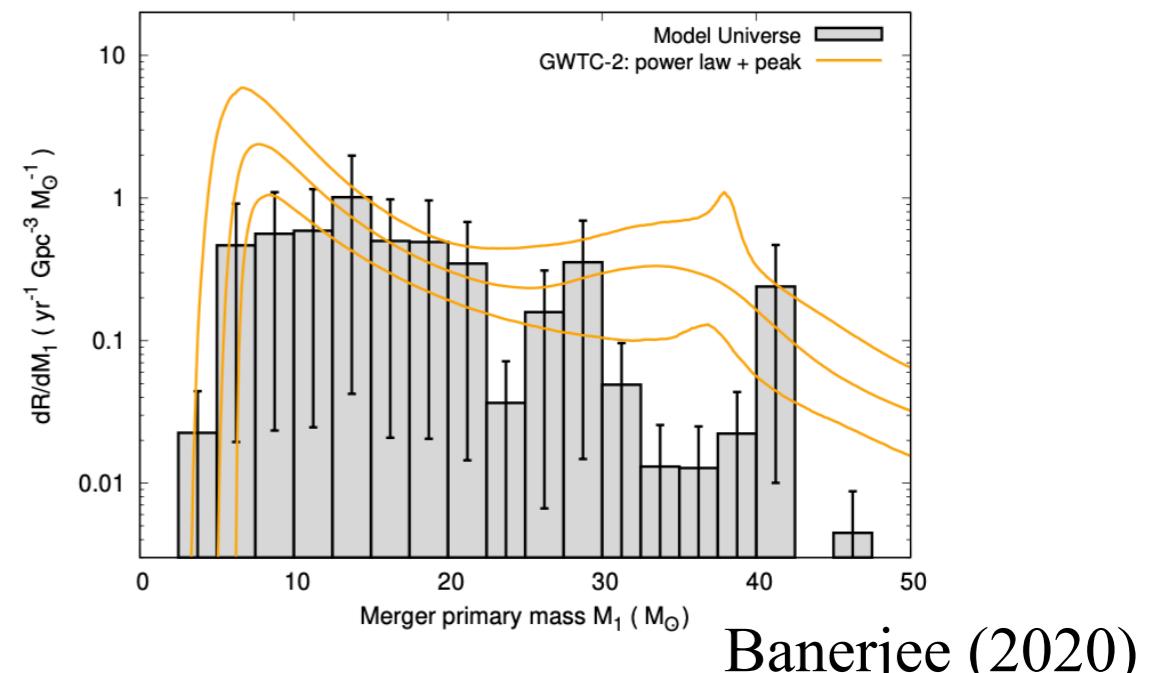
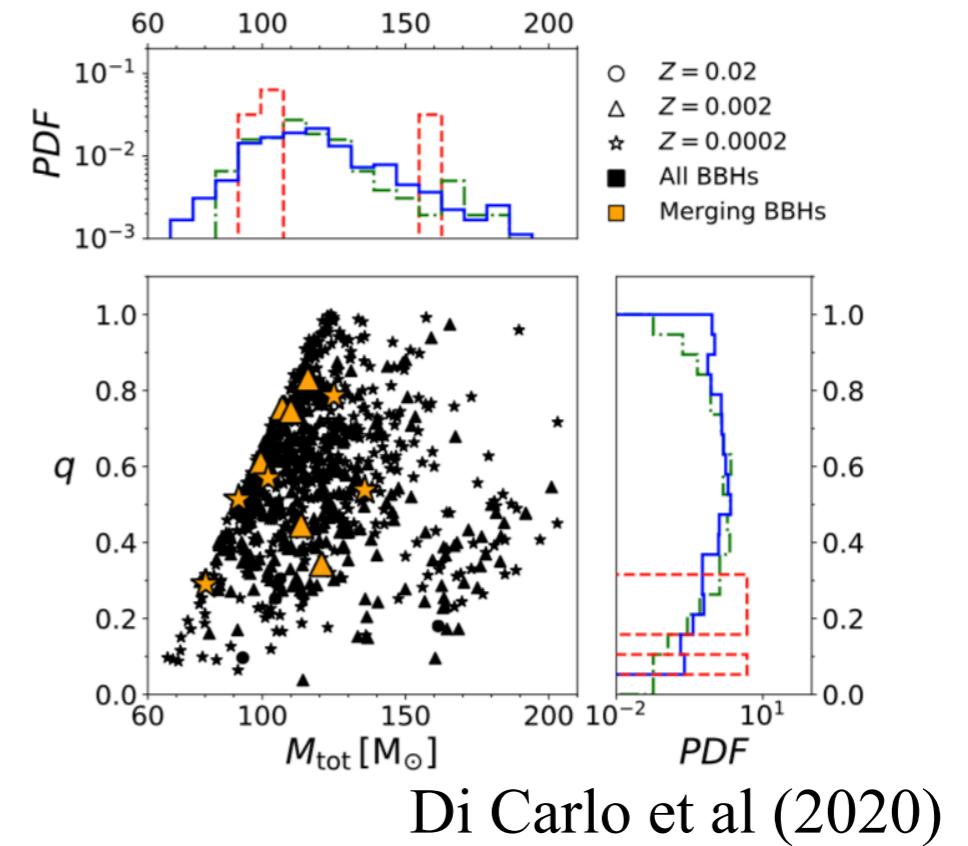


Open clusters

- Enough BH budget
- Common envelope and dynamical capture (Kumamoto et al. 2019; 2020)
- Mass gap BHs (Di Carlo et al. 2019; 2020) via post-MS and MS mergers (Kirihara's study can apply to)
- Strong PPI/PISN model (No $> 40M_{\odot}$ BH without stellar and BH mergers)

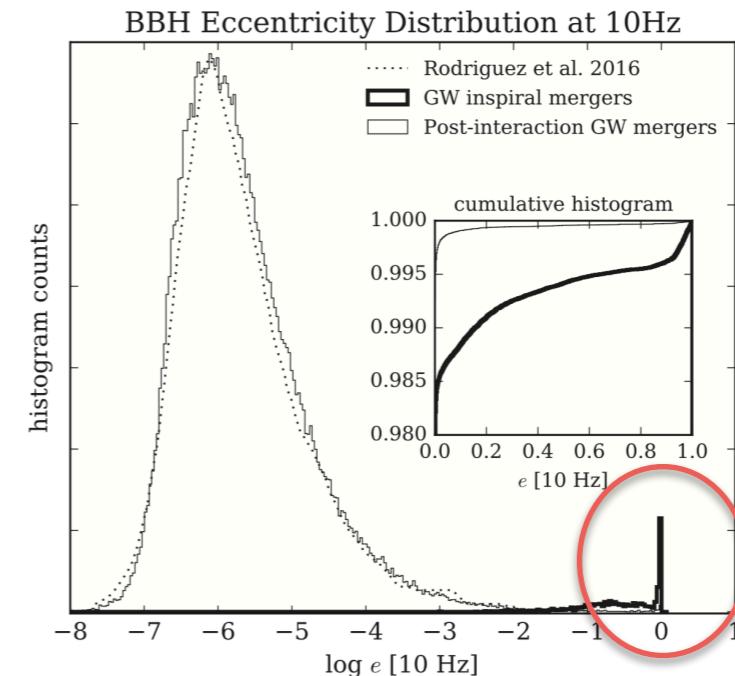


Kumamoto et al. (2019;2020)



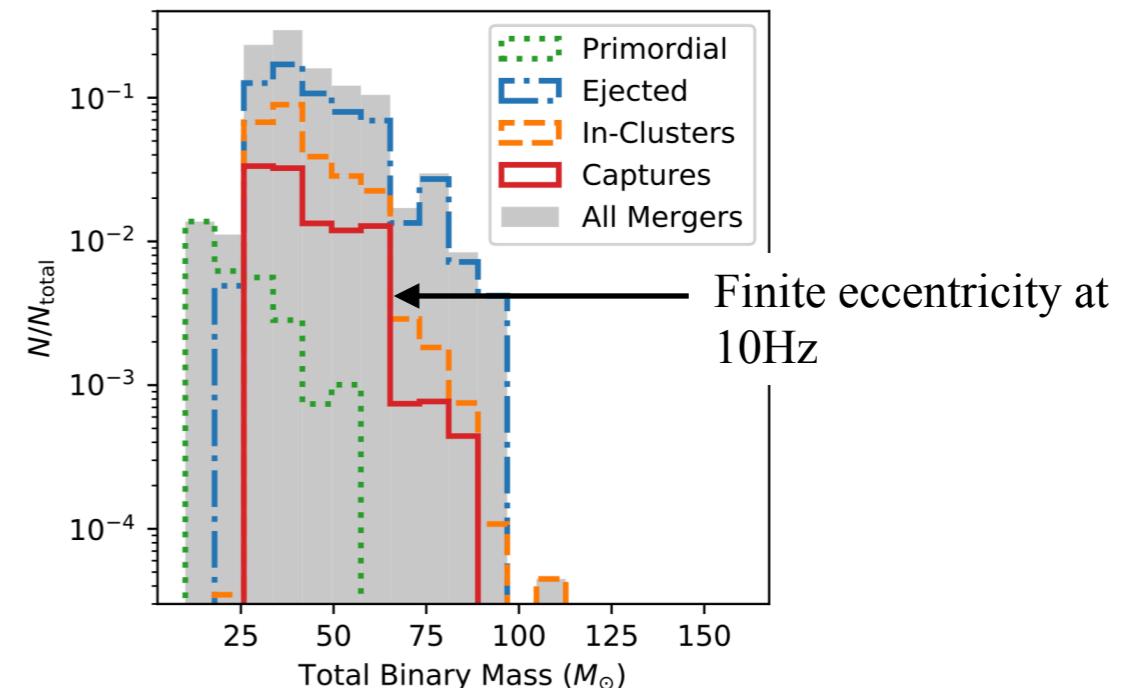
BH-BH eccentricity

- ~1% of BH-BH mergers leave finite eccentricities at LIGO/Virgo/KAGRA bands.
- Their masses are similar to those of in-cluster mergers, but smaller than those of ejected BH-BHs.
- The present-day BHs in GCs are lighter than previously ejected BHs.
- GW190521 **might** be an eccentric merger (Gayathri et al. 2020; Romero-Shaw et al. 2020).
 - But, people who believe it are supporters of the AGN scenario (Samsing et al. 2020; Tagawa et al. 2020).



Samsing, Ramirez-Ruiz (2017)

1st-Generation BBH Mergers, $z < 1$



Rodriguez et al. (2018)

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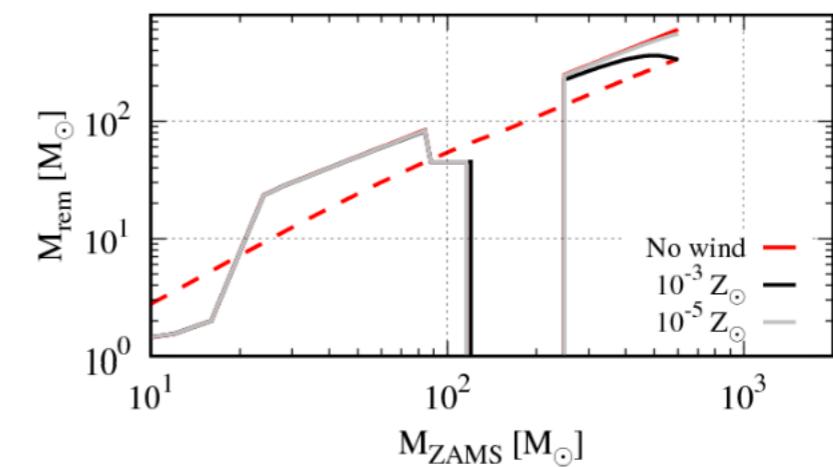
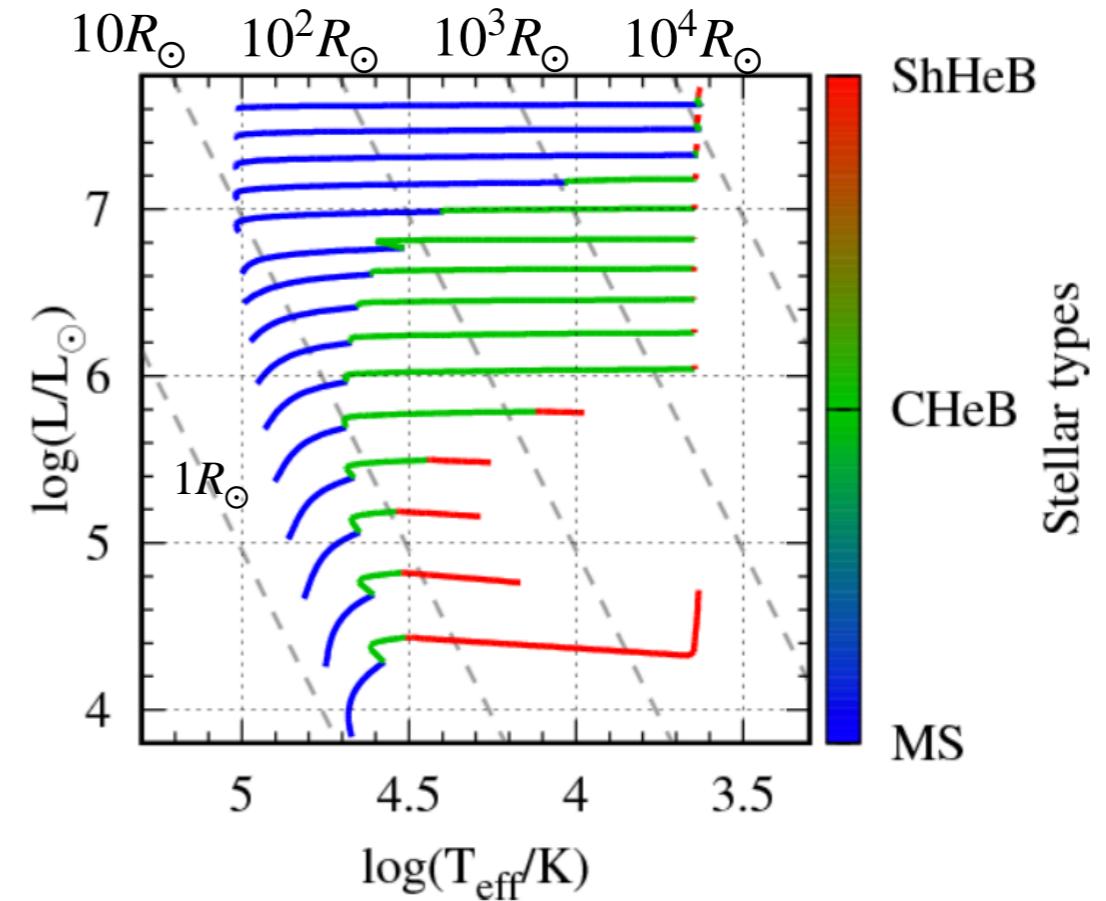
- From O1/O2 to O3a
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 - Tanikawa, Yoshida, Kinugawa, Takahashi, Umeda (2020a, MNRAS, 495, 4170)
 - Tanikawa, Susa, Yoshida, Trani, Kinugawa (2020b, arXiv:2008.01890)
 - Tanikawa, Kinugawa, Yoshida, Hijikawa, Umeda (2020c, arXiv:2010.07616)

Black hole budget (Pop. III)

- BH-BH density: $n_{\text{BH-BH}} \sim 10^{11} \text{Gpc}^{-3} \left(\frac{\Gamma_{\text{BH-BH}}}{10 \text{Gpc}^{-3} \text{yr}^{-1}} \right) \left(\frac{T_{\text{Hubble}}}{10 \text{Gyr}} \right)$
 - Binary evolution:
 - $n_{\text{BH,star}} \sim 10^{15} \text{Gpc}^{-3} \left(\frac{\rho_{\text{star}}}{10^{18} M_{\odot} \text{Gpc}^{-3}} \right) \left(\frac{\eta_{\text{BH}}}{10^{-3} M_{\odot}} \right)$ Abbott et al. (2020)
 - $\frac{n_{\text{BH-BH}}}{n_{\text{BH,star}}} \sim 10^{-4} \left(\frac{\Gamma_{\text{BH-BH}}}{10 \text{Gpc}^{-3} \text{yr}^{-1}} \right)$ Kroupa (2001), $M_{\text{zams}} \gtrsim 20 M_{\odot} \rightarrow \text{BH}$
 - Pop. III binary evolution
 - $n_{\text{BH,III}} \sim 10^{11} \text{Gpc}^{-3} \left(\frac{\rho_{\text{star}}}{10^{13} M_{\odot} \text{Gpc}^{-3}} \right) \left(\frac{\eta_{\text{BH}}}{10^{-2} M_{\odot}} \right)$ Madau, Dickinson (2014)
 - $\frac{n_{\text{BH-BH}}}{n_{\text{BH,III}}} \sim 1 \left(\frac{\Gamma_{\text{BH-BH}}}{10 \text{Gpc}^{-3} \text{yr}^{-1}} \right)$
 - $f(M) \propto M^{-1}$ ($10 \leq M/M_{\odot} \leq 100$)
 - $M_{\text{zams}} \gtrsim 20 M_{\odot} \rightarrow \text{BH}$
- Magg et al. (2016); Skinner, Wise (2020);
but de Souza et al. (2011); Inayoshi et al. (2016)

Binary population synthesis

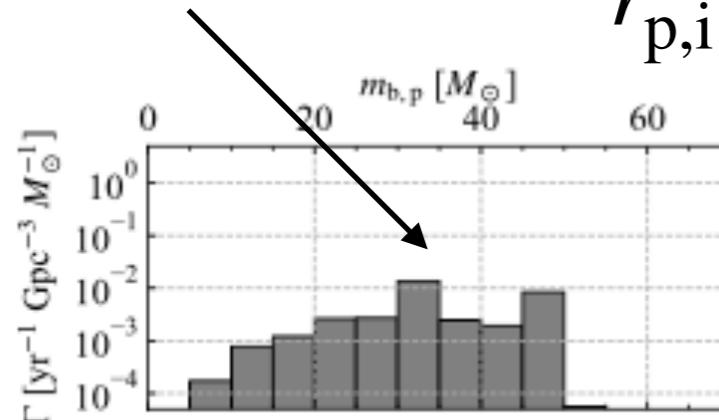
- $10 \leq m_{1,\text{zams}}/M_\odot \leq 300$
- $10/m_{1,\text{zams}} \leq m_{2,\text{zams}}/m_{1,\text{zams}} \leq 1$
- $r_{\text{p,i}} \geq 10$ or $200R_\odot$
- Pop. III model with large convective overshooting (L model)
- Mass transfer, common envelope etc.
- Fryer's SN model with PPI/PISN
- One Pop. III binary per minihalo



Tanikawa et al. (2020a)

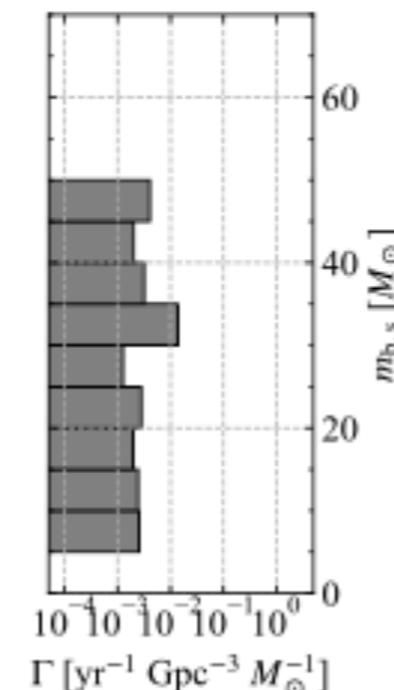
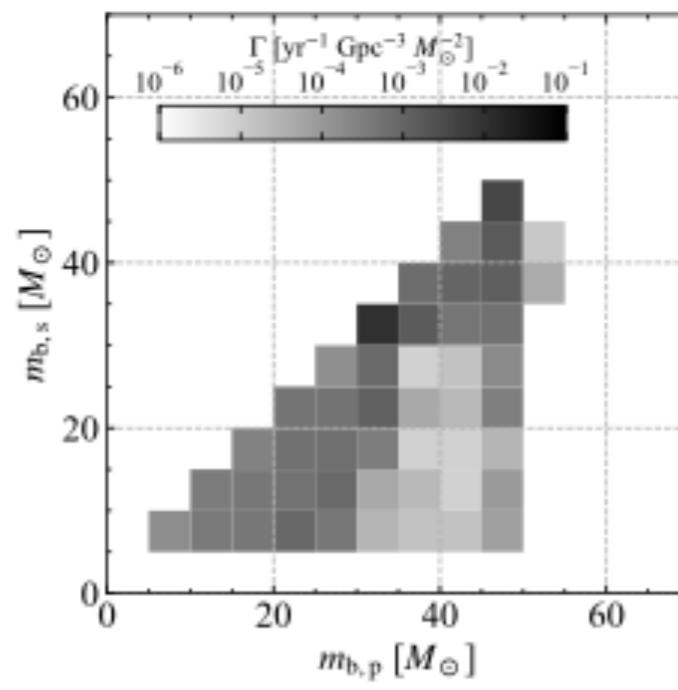
Pop. III BH-BHs

Kinugawa's peak

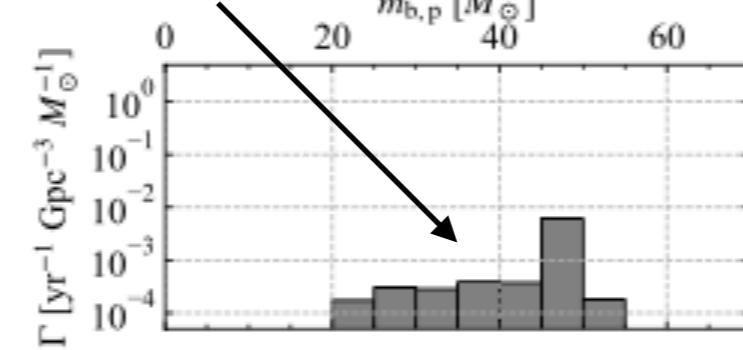


$$r_{\text{p,i}} \geq 10R_{\odot}$$

MT path
CE path

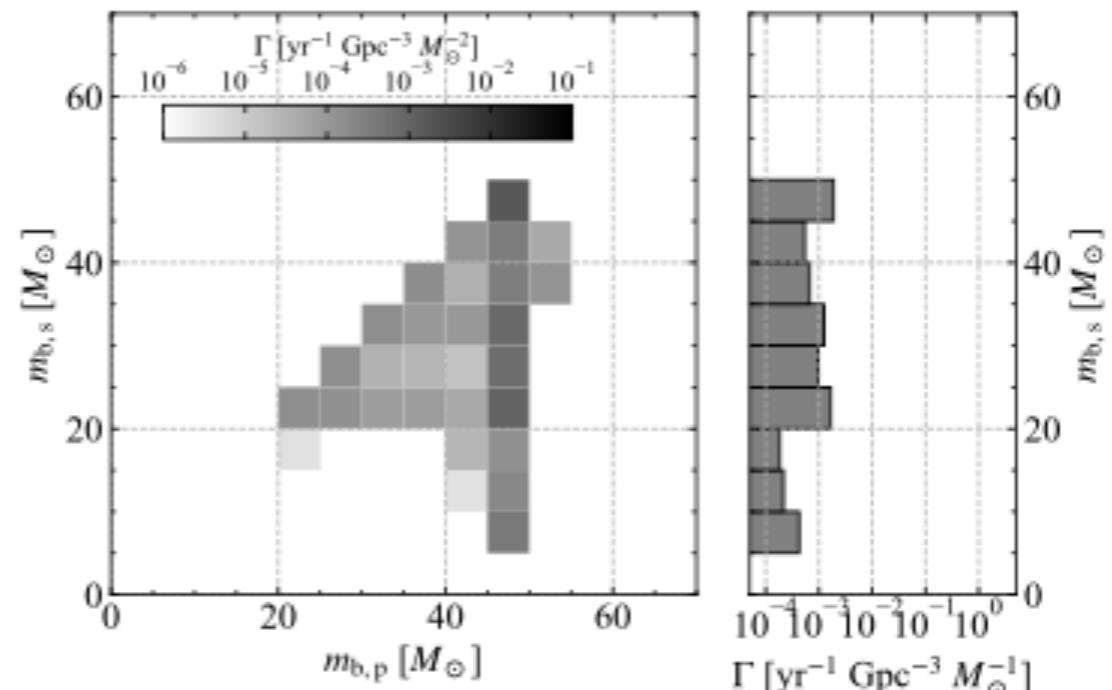


w/o Kinugawa's peak



$$r_{\text{p,i}} \geq 200R_{\odot}$$

ratio 0.0a2e2
CE path

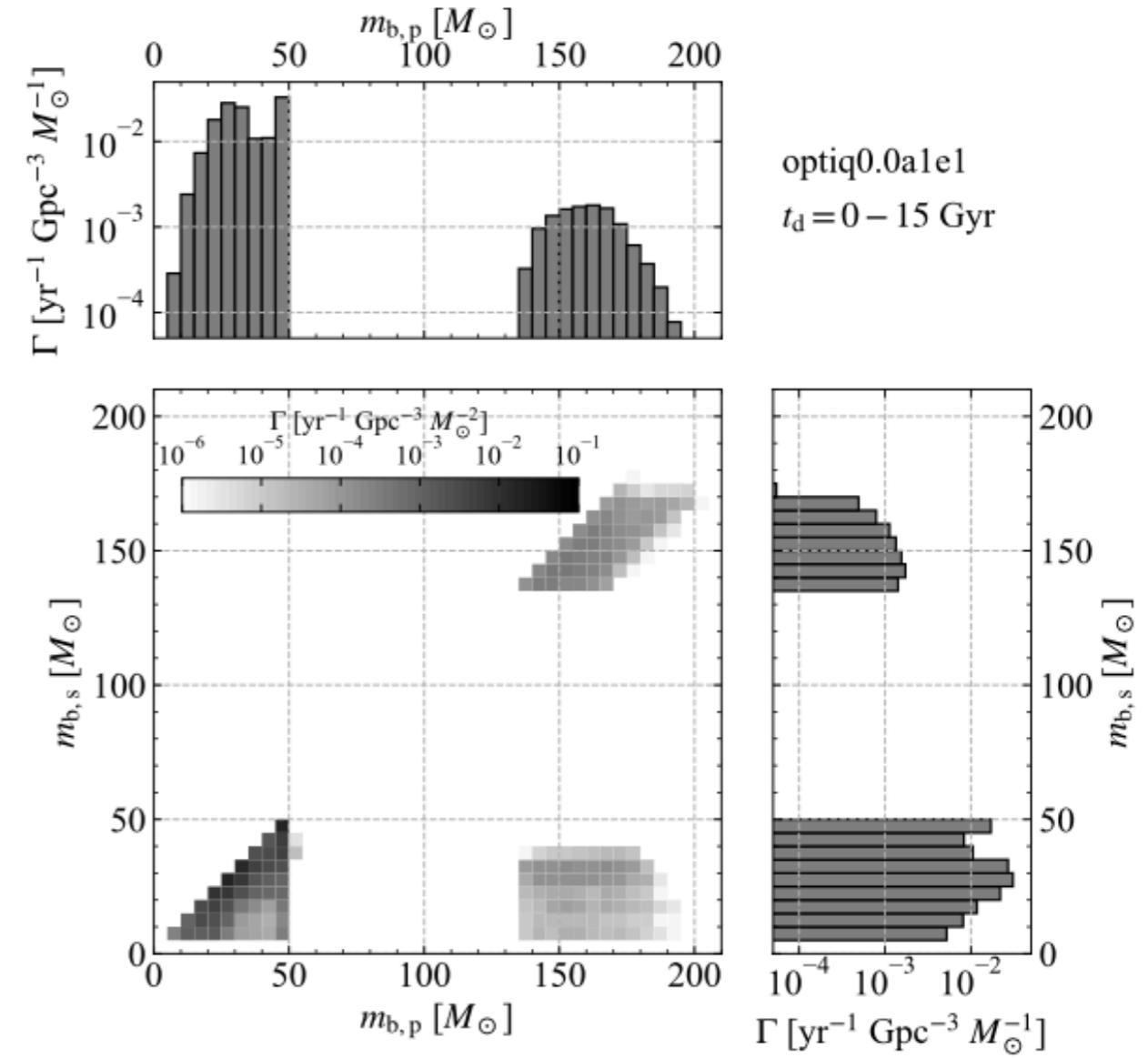
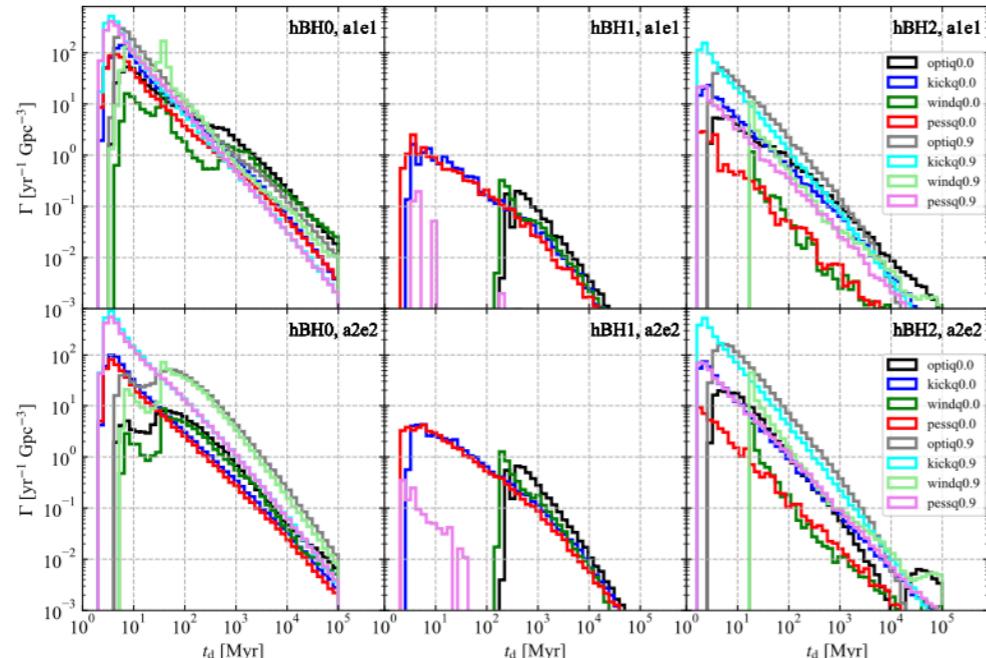


$$\Gamma_{m_1 \lesssim 50M_{\odot}} \sim 0.1 \text{ Gpc}^{-3} \text{ yr}^{-1}$$

Tanikawa et al. (2020b)

Pop. III BH-BHs

- $\Gamma_{m_1 \lesssim 50M_\odot} \sim 0.1 \text{ Gpc}^{-3} \text{ yr}^{-1}$
- $\Gamma_{50M_\odot \lesssim m_1 \lesssim 130M_\odot} \sim 0 \text{ Gpc}^{-3} \text{ yr}^{-1}$
- $\Gamma_{m_1 \gtrsim 130M_\odot} \sim 0.01 \text{ Gpc}^{-3} \text{ yr}^{-1}$
 - GW190521 can be $m_1 \gtrsim 130M_\odot$ (e.g. Fishbach, Holz 2020).



Tanikawa et al. (2020b)

GW190521

- Merger of $85^{+21}_{-14} M_{\odot}$ and $66^{+17}_{-18} M_{\odot}$ BHs
- The primary BH has only a 0.32% probability of being below $65 M_{\odot}$.
- At least one BH lies within the PI mass gap.

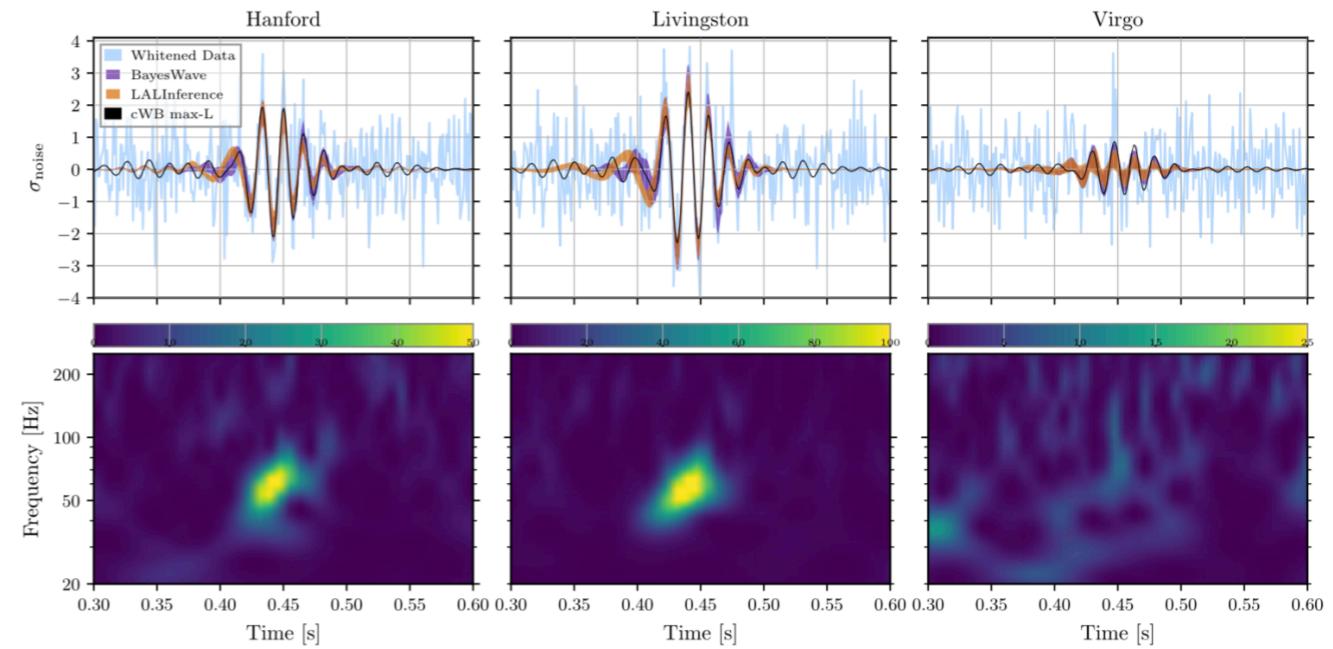
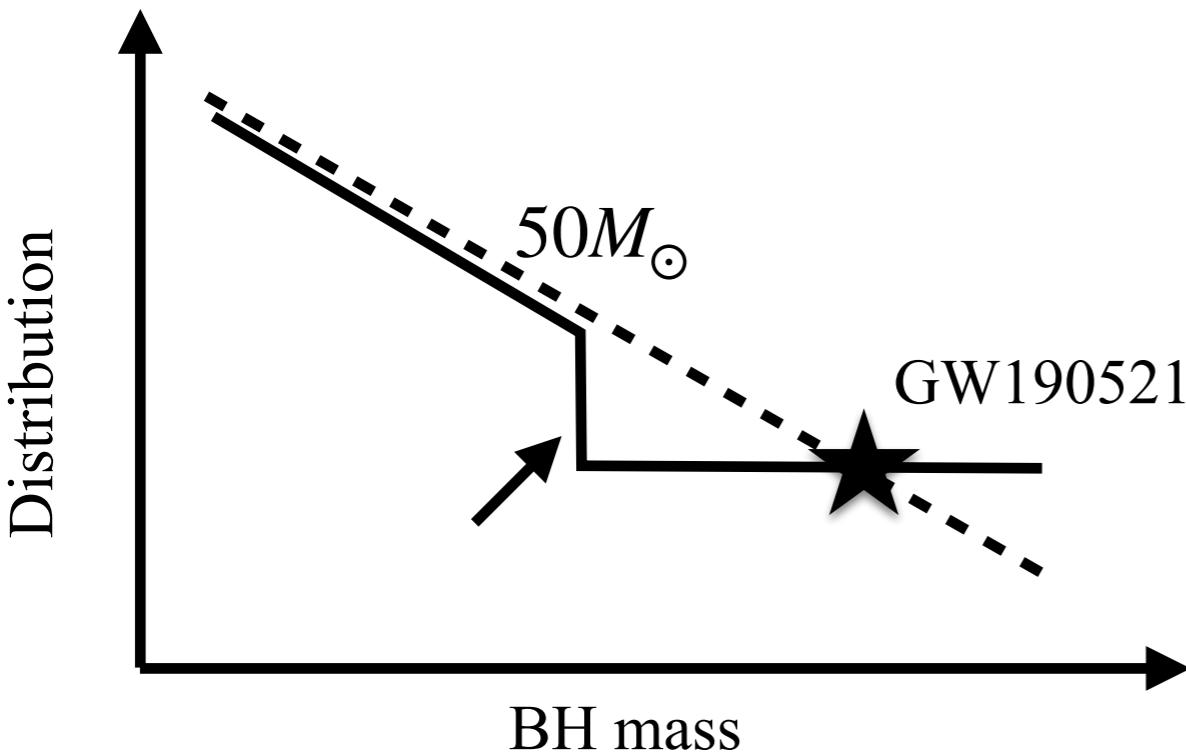


TABLE I. Parameters of GW190521 according to the NRSur7dq4 waveform model. We quote median values with 90% credible intervals that include statistical errors.

Parameter	
Primary mass	$85^{+21}_{-14} M_{\odot}$
Secondary mass	$66^{+17}_{-18} M_{\odot}$
Primary spin magnitude	$0.69^{+0.27}_{-0.62}$
Secondary spin magnitude	$0.73^{+0.24}_{-0.64}$
Total mass	$150^{+29}_{-17} M_{\odot}$
Mass ratio ($m_2/m_1 \leq 1$)	$0.79^{+0.19}_{-0.29}$
Effective inspiral spin parameter (χ_{eff})	$0.08^{+0.27}_{-0.36}$
Effective precession spin parameter (χ_p)	$0.68^{+0.25}_{-0.37}$
Luminosity Distance	$5.3^{+2.4}_{-2.6} \text{ Gpc}$
Redshift	$0.82^{+0.28}_{-0.34}$
Final mass	$142^{+28}_{-16} M_{\odot}$
Final spin	$0.72^{+0.09}_{-0.12}$
$P(m_1 < 65 M_{\odot})$	0.32%
\log_{10} Bayes factor for orbital precession	$1.06^{+0.06}_{-0.06}$
\log_{10} Bayes factor for nonzero spins	$0.92^{+0.06}_{-0.06}$
\log_{10} Bayes factor for higher harmonics	$-0.38^{+0.06}_{-0.06}$

Abbott et al. (2020)

Single star evolution

- It is not hard to form mass-gap BH through single star evolution.
- Formation Process
 - A star with $M_{\text{zams}} \sim 90M_{\odot}$ and $Z \sim 0.01Z_{\odot}$.
 - Evolution to a BH progenitor with $M_{\text{tot}} \sim 90M_{\odot}$ and $M_{\text{c,He}} \lesssim 40M_{\odot}$.
 - Collapse to $\sim 90M_{\odot}$ BH without PPI/PISN owing to small He core mass.
 - Light He core, massive H envelope

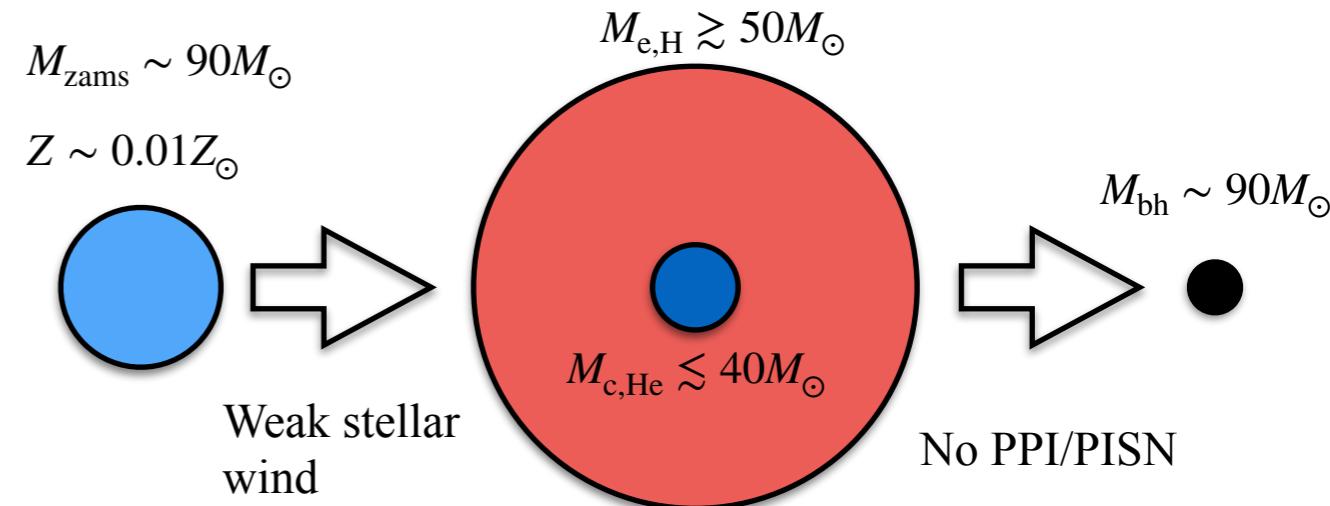


Table 1. Summary of the results

M_{ini} (M_{\odot})	M_{CO} (M_{\odot})	M_{He} (M_{\odot})	# of PPI	Ejection #	Log T_{peak} (K)	M_{rem} (M_{\odot})	Ejecta Energy (10^{50} erg)
L Models ($f_{\text{OV}}=0.03$)							
70	34.2	38.9-48.8	4	0	-	70	-
75	34.9	39.3	4	1	9.81	42.4	6.5
80	37.4	42.2-42.9	3	1	9.71	42.4	0.18
100	48.1	53.6	2	1	9.65	52.2	4.5
120	57.9	64.9	1	1	9.66	60.3	4.7
135	65.4	73.5	1	1	9.63	66.9	5.6
M Models ($f_{\text{OV}}=0.01$)							
70	27.0	30.3-34.4	0	0	-	70	-
80	31.8	35.3-39.4	5	0	-	80	-
90	37.2	41.9-44.8	3	1	9.76	83.0	1.4
100	42.7	47.2-50.1	0	1	9.72	61.7	1.2

Pop. III

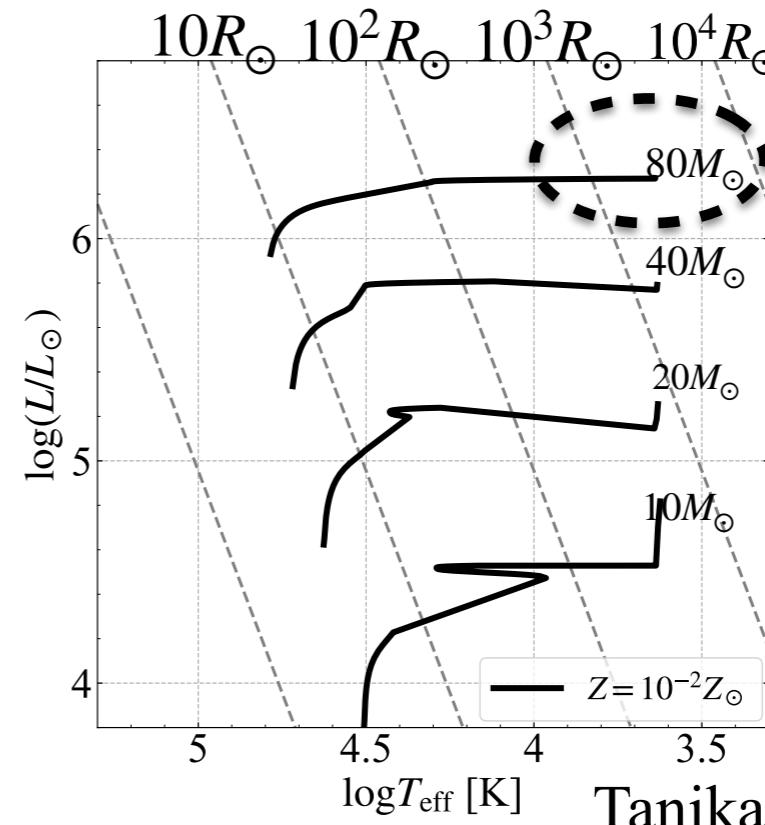
$$M_{\text{zams}} = 70M_{\odot} \rightarrow M_{\text{bh}} = 70M_{\odot}$$

$$M_{\text{zams}} = 80M_{\odot} \rightarrow M_{\text{bh}} = 80M_{\odot}$$

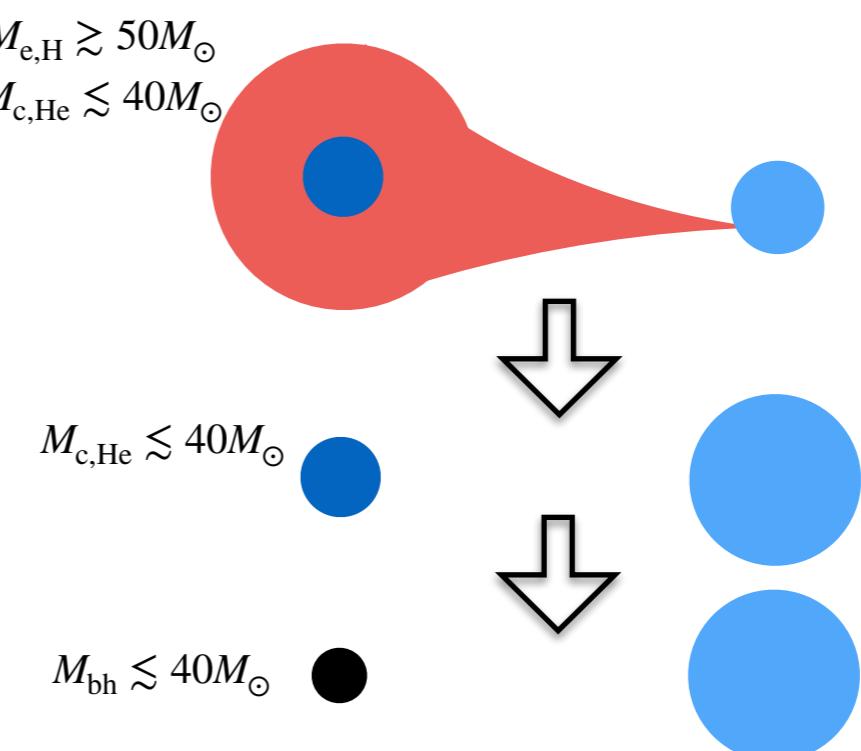
	(interval: 3.1 yr)		2	9.80	77.7	39	
35	58.3	65.3-72.6	1	1	9.64	108.7	5.6

Binary star evolution

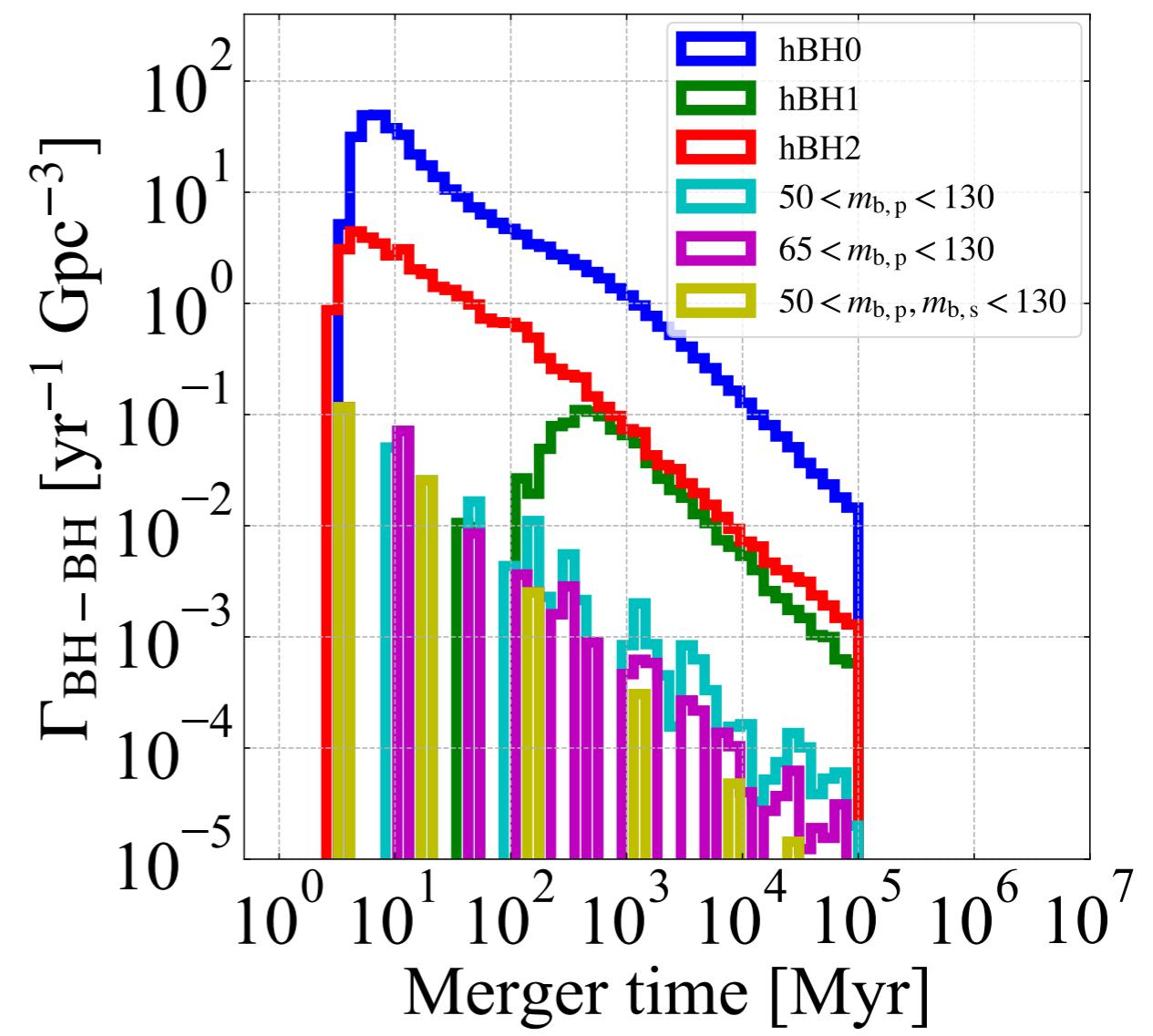
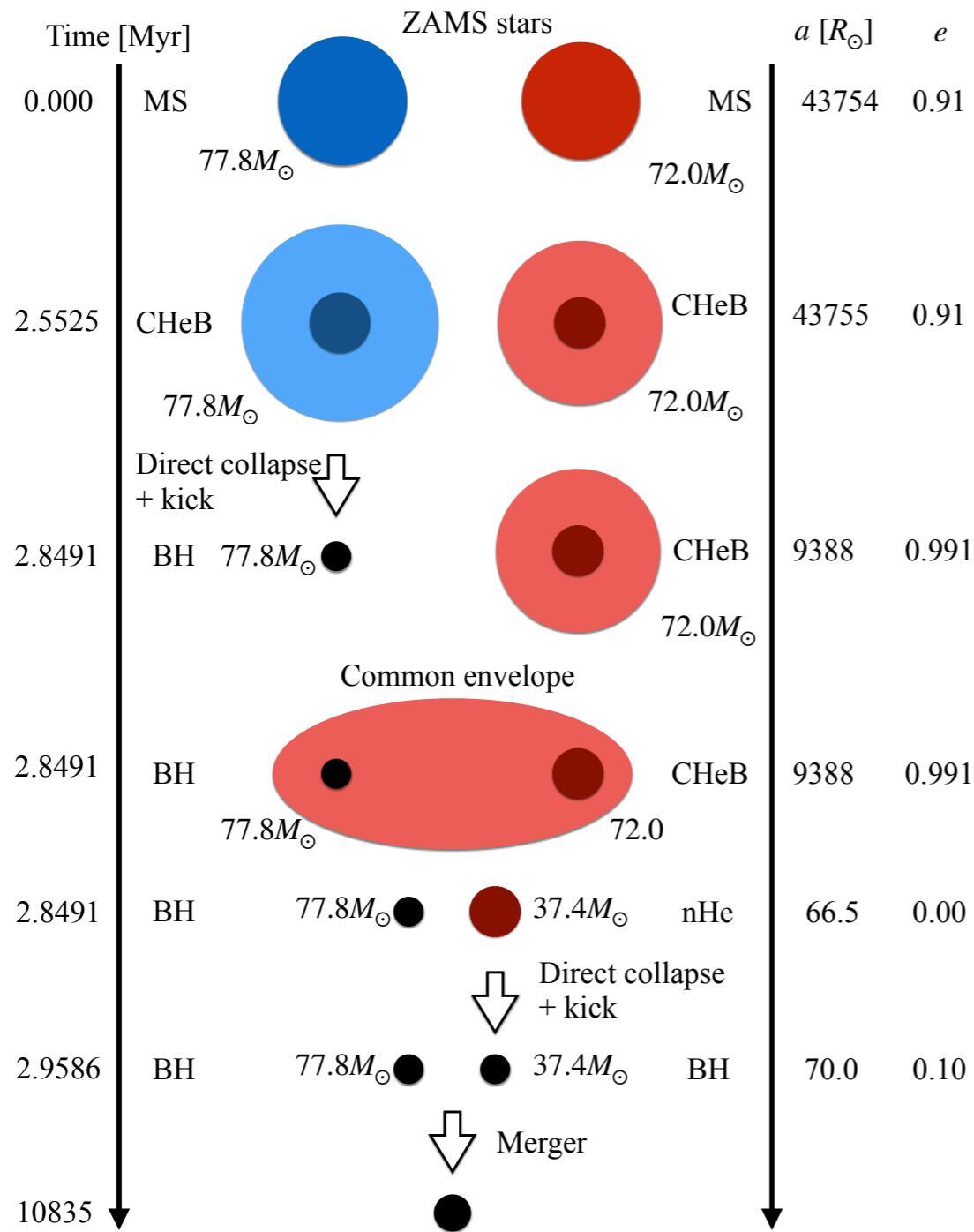
- Merger of $85M_{\odot}$ and $66M_{\odot}$ BHs
 - Merger time $\lesssim 10$ Gyr
 - $a \lesssim 10^2R_{\odot}$, $e \sim 0$
- A star with $M_{\text{zams}} \gtrsim 80M_{\odot}$ expands to $R \gtrsim 10^3R_{\odot}$.
- The star loses its H envelope, stripped by its companion star.
- No massive H envelope, no mass-gap BH.



Tanikawa et al. (2020a)



BH natal kick



Tanikawa et al. (2020b)

Summary

- Globular clusters can fill the pair instability mass gap.
 - But, the event rate ratio does not seem consistent.
- Pop. III merger rate with $m_1 \gtrsim 130M_\odot$ is $\sim 0.01 \text{ Gpc}^{-3} \text{ yr}^{-1}$.
 - It may be consistent with GW190521 if GW190521 contains $m_1 \gtrsim 130M_\odot$ (e.g. Fishbach, Holtz 2020)
- Whether Pop. III binaries can form GW190521 strongly depends on convective overshooting.